# HALF TWISTS OF HODGE STRUCTURES OF CM-TYPE 

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#### Abstract

To a Hodge structure $V$ of weight $k$ with CM by a field $K$ we associate Hodge structures $V_{-n / 2}$ of weight $k+n$ for $n$ positive and, under certain circumstances, also for $n$ negative. We show that these 'half twists' come up naturally in the Kuga-Satake varieties of weight two Hodge structures with CM by an imaginary quadratic field.


## Introduction

A Hodge structure of CM-type is a Hodge structure $V$ on which a CM-field $K$ acts. Given a CM-type (a set of certain complex embeddings of $K$ ), we give a simple construction of Hodge structures $V_{-n / 2}$ of weight $k+n$, for any positive integer $n$ where $k$ is the weight of $V$. These half twists of $V$ are not related to the Tate twists $V(-n)$ of $V: V_{-n} \neq V(-n)$.

In certain circumstances one can also define the half twist $V_{1 / 2}$, of weight $k-1$. The geometry underlying the half twist in the case of hypersurfaces in projective space is investigated in [VGI]. A basic case is Example 2.12 in this paper.

In the first section we recall the basic definitions of Hodge structures and we define the half twist. In the second section we look at half twists from the point of view of representations of $\mathbf{C}^{*}$. We also give a geometrical interpretation of the half twist in case $V$ is a sub-Hodge structure of $H^{k}(X, \mathbf{Q})$ for a smooth projective variety $X$. In the last section we consider Hodge structures $V$ of weight two with $\operatorname{dim} V^{2,0}=1$ and with a suitable action of an imaginary quadratic field. We show that the first cohomology group of the Kuga-Satake variety associated to such a Hodge structure has a summand which is the half twist of $V$ and that half twists can be used to understand the other summands as well. So, in a certain sense, half twists 'partly generalize' the Kuga-Satake construction which associates weight one Hodge structures to certain weight two Hodge structures. These results were motivated by an example of C. Voisin [V] which also inspired the general definition of half twist.

I'm indebted to E. Izadi and M. Nori for simplifications and generalizations of my original construction.

## 1. Definitions and basic properties

1.1. Hodge structures. Recall that a (rational) Hodge structure of weight $k\left(\in \mathbf{Z}_{\geq 0}\right)$ is a $\mathbf{Q}$-vector space $V$ with a decomposition of its complexification $V_{\mathbf{C}}:=V \otimes_{\mathbf{Q}} \mathbf{C}$ (where complex conjugation is given by $\overline{v \otimes z}:=v \otimes \bar{z}$ for $v \in V$ and $z \in \mathbf{C})$ :

$$
V_{\mathbf{C}}=\oplus_{p+q=k} V^{p, q}, \quad \text { and } \quad \overline{V^{p, q}}=V^{q, p}, \quad\left(p, q \in \mathbf{Z}_{\geq 0}\right)
$$

Note that we insist on $p$ and $q$ being non-negative integers throughout this paper, so we only consider 'effective' Hodge structures.
1.2. Hodge structures of CM-type. Let $V$ be a Hodge structure such that $V$ is also a $K$-vector space for some $C M$-field $K$ and such that the Hodge decomposition on $V$ is stable under the action of $K$ :

$$
x V^{p, q} \subset V^{p, q}, \quad\left(x \in K, p, q \in \mathbf{Z}_{\geq 0}\right)
$$

In particular, $K \hookrightarrow E n d d_{H o d}(V)$. We will then say that $V$ is a Hodge structure of CM-type (with field $K$ ).
1.3. CM-types. To define the half twist we need to fix a CM-type of the field $K$. Recall that the CM-field $K$ has $2 r=[K: \mathbf{Q}]$ complex embeddings $K \hookrightarrow \mathbf{C}$ and that a CM-type is a subset $\Sigma=\left\{\sigma_{1}, \ldots, \sigma_{r}\right\}$ of distinct embeddings with the property that no two are complex conjugate. Hence if we define as us ual $\bar{\Sigma}:=\left\{\bar{\sigma}_{1}, \ldots, \bar{\sigma}_{r}\right\}$ then any embedding of $K$ is either in $\Sigma$ or in $\bar{\Sigma}$.
1.4. The half twists. Let $V$ be a Hodge structure of CM-type with field $K$. We consider the eigenspaces of the $K$-action on the $V^{p, q}$ 's:

$$
V_{\sigma}^{p, q}:=\left\{v \in V^{p, q}: x v=\sigma(x) v \quad \forall x \in K\right\}, \quad \sigma: K \hookrightarrow \mathbf{C} .
$$

Given a CM-type $\Sigma$, we define two subspaces of $V^{p, q}$ whose direct sum is $V^{p, q}$ :

$$
V_{\Sigma}^{p, q}:=\oplus_{\sigma \in \Sigma} V_{\sigma}^{p, q}, \quad V_{\bar{\Sigma}}^{p, q}:=\oplus_{\sigma \in \bar{\Sigma}} V_{\sigma}^{p, q}
$$

The negative half twist of $V$, denoted by $V_{-1 / 2}\left(=V_{\Sigma,-1 / 2}\right)$, is the decomposition of $V_{\mathbf{C}}$ given by the subspaces:

$$
V_{-1 / 2}^{r, s}:=V_{\Sigma}^{r-1, s} \oplus V_{\bar{\Sigma}}^{r, s-1} .
$$

It is not hard to see that this is a Hodge structure, of CM-type with field $K$, on $V$ of weight $k+1$ where $k$ is the weight of $V$. By successively performing the negative half twist one obtains $V_{-n / 2}$, a Hodge structure on $V$ of weight $k+n$.

We observe that to define the (positive) half twist one would put:

$$
V_{1 / 2}^{r, s}:=V_{\Sigma}^{r+1, s} \oplus V_{\bar{\Sigma}}^{r, s+1}
$$

however, now the subspaces $V_{\bar{\Sigma}}^{k, 0}$ and $V_{\Sigma}^{0, k}$ of $V_{\mathbf{C}}$ do not appear in $\left(V_{1 / 2}\right)_{\mathbf{C}}$ and therefore this definition does not define a Hodge structure on $V$ (and in general not on any Q-subspace of $V)$. In particular, we will define the half twist of $V$ only if $V_{\bar{\Sigma}}^{k, 0}=0$ (the complex conjugate of this space is $V_{\Sigma}^{0, k}$ which is then also trivial).

## 2. Half twists via representations

2.1. Hodge structures can also be defined via representations of $\mathbf{C}^{*}$ (more precisely, algebraic representations of $\operatorname{Res}_{\mathbf{C} / \mathbf{R}}\left(G_{m}\right)$ in $G L\left(V_{\mathbf{R}}\right)$ ). We determine the representations corresponding to the half twists. Proposition 2.8 often points out interesting geometry, as in example 2.12. In VGI more such examples are investigated.
2.2. The Hodge structure on a $\mathbf{Q}$-vector space $V$ defined by an algebraic representation $h$ : $\mathbf{C}^{*} \rightarrow G L\left(V_{\mathbf{R}}\right)$ will be denoted by $(V, h)$, its Hodge decomposition is:

$$
V^{p, q}:=\left\{v \in V_{\mathbf{C}}: h(z) v=z^{p} \bar{z}^{q} v\right\} .
$$

The usual algebra constructions on representations can be applied to Hodge structures. In particular, given rational Hodge structures $(V, h),\left(W, h_{W}\right)$ of weight $k, k_{W}$ their tensor product is the rational Hodge structure of weight $k+k_{W}$ defined by:

$$
h \otimes h_{W}: \mathbf{C}^{*} \longrightarrow G L\left(V_{\mathbf{R}} \otimes W_{\mathbf{R}}\right), \quad z \longmapsto\left[v \otimes w \mapsto(h(z) w) \otimes\left(h_{W}(z) w\right)\right] .
$$

2.3. Let $(V, h)$ be a Hodge structure of weight $k$ of CM-type with field $K$. To identify the $\mathbf{C}^{*}$-representation on $V_{\mathbf{R}}$ which defines the half twist, consider the $\mathbf{R}$-linear extension of the action of $K$ on $V_{\mathbf{R}}$. This gives an action of $K \otimes_{\mathbf{Q}} \mathbf{R}$ on $V_{\mathbf{R}}$ and recall that $K \otimes_{\mathbf{Q}} \mathbf{R} \cong \oplus_{i=1}^{r} \mathbf{C}$ (if you write

$$
K=\mathbf{Q}[X] /(f) \quad \text { then } \quad K \otimes \mathbf{R} \cong \mathbf{R}[X] /(f) \cong \prod_{i=1}^{r} \mathbf{R}[X] /\left(f_{i}\right)
$$

where $f=\Pi f_{i}$ is the decomposition of $f$ in irreducible polynomials $f_{i}$ of degree 2 in $\mathbf{R}[X]$ ). Choosing a CM-type $\Sigma=\left\{\sigma_{1}, \ldots, \sigma_{r}\right\}$ of $K$ specifies an isomorphism

$$
K \otimes \mathbf{R} \xrightarrow{\cong} \oplus_{j=1}^{r} \mathbf{C}, \quad x \otimes t \longmapsto\left(t \sigma_{1}(x), \ldots, t \sigma_{r}(x)\right),
$$

we will denote the $\mathbf{R}$-linear extensions of the $\sigma_{i}$ 's by the same symbols. The inverse of this isomorphism is denoted by $\phi$

$$
\phi: \oplus_{j=1}^{r} \mathbf{C} \xrightarrow{\cong} K \otimes \mathbf{R}, \quad \text { with } \quad \sigma_{i}\left(\phi\left(z_{1}, \ldots, z_{r}\right)\right)=z_{i} \quad(1 \leq i \leq r) .
$$

Denoting by $x \mapsto \bar{x}$ the complex conjugation on $K$ as well as its $\mathbf{R}$-linear extension to $K \otimes_{\mathbf{Q}} \mathbf{R}$ let $\sigma_{r+i}(x):=\overline{\sigma_{i}(x)}=\sigma_{i}(\bar{x})$, hence

$$
\sigma_{r+i}\left(\phi\left(z_{1}, \ldots, z_{r}\right)\right)=\bar{z}_{i} \quad(1 \leq i \leq r) .
$$

We define a homomorphism (depending on the CM-type $\Sigma$ ) by composing the diagonal inclusion $\Delta$ with the isomorphism $\phi=\phi_{\Sigma}$ :

$$
g_{\Sigma}: \mathbf{C}^{*} \xrightarrow{\Delta} \prod_{i=1}^{r} \mathbf{C}^{*} \xrightarrow{\phi}\left(K \otimes_{\mathbf{Q}} \mathbf{R}\right)^{*}
$$

Since $\left(K \otimes_{\mathbf{Q}} \mathbf{R}\right)^{*}$ acts on $V_{\mathbf{R}}$, we get a representation, also denoted by $g_{\Sigma}$ of $\mathbf{C}^{*}$ on $V_{\mathbf{R}}$. As $V$ is of CM-type, the actions of $K$ and $h\left(\mathbf{C}^{*}\right)$ commute and thus the action of $g_{\Sigma}\left(\mathbf{C}^{*}\right)\left(\subset\left(K \otimes_{\mathbf{Q}} \mathbf{R}\right)^{*}\right)$ on $V_{\mathbf{R}}$ commutes with the action of $h\left(\mathbf{C}^{*}\right)$.

To understand the action of $g_{\Sigma}$ we observe that the eigenvalues of $x \in K$ on $V_{\mathbf{C}}$ are the $\sigma_{i}(x)$, $1 \leq i \leq 2 r$, each with the same multiplicity. The decomposition in eigenspaces is

$$
V_{\mathbf{C}}=\oplus_{\sigma \in \Sigma \cup \bar{\Sigma}} V_{\mathbf{C}, \sigma} \quad \text { with } \quad x v=\sigma(x) v \quad\left(x \in K \otimes_{\mathbf{Q}} \mathbf{R}, v \in V_{\mathbf{C}, \sigma}\right)
$$

In particular, if $x=g_{\Sigma}(z)$ then since $g_{\Sigma}(z)=\phi(\Delta(z))=\phi(z, \ldots, z)$ we get:

$$
g_{\Sigma}(z) v=\phi(z, \ldots, z) v= \begin{cases}z v & v \in \oplus_{\sigma \in \Sigma} V_{\mathbf{C}, \sigma} \\ \bar{z} v & v \in \oplus_{\sigma \in \bar{\Sigma}} V_{\mathbf{C}, \sigma}\end{cases}
$$

This leads to an alternative, but equivalent, definition of the half twist:
2.4. Definition. Let $(V, h, K)$ be a Hodge structure of CM-type and let $\Sigma$ be a CM-type of $K$. For $n \in \mathbf{Z}$, the $n$-th half twist $V_{-n / 2}$ is the $\mathbf{C}^{*}$-representation defined by the homomorphism

$$
g_{\Sigma}^{n} h: \mathbf{C}^{*} \longrightarrow G L\left(V_{\mathbf{R}}\right), \quad z \longmapsto g_{\Sigma}^{n}(z) h(z) .
$$

2.5. For positive $n$ the representation space $V_{-n / 2}$ is a Hodge structure of weight $k+n$ where $k$ is the weight of $V$. Moreover, with the same CM-type, one has:

$$
\left(V_{n / 2}\right)_{m / 2}=V_{(n+m) / 2}
$$

We already observed that the half twist $V_{1 / 2}$ is a Hodge structure iff the eigenspaces for the $K$-action $V_{\sigma}^{k, 0}$ are trivial for $\sigma \in \bar{\Sigma}$.
2.6. Tate twists. The Tate Hodge structure $\mathbf{Q}(n)(n \in \mathbf{Z})$ is defined by the vector space $\mathbf{Q}$ and the homomorphism:

$$
h_{n}: \mathbf{C}^{*} \longrightarrow G L_{1}(\mathbf{R}), \quad z \longmapsto(z \bar{z})^{-n},
$$

it has weight $-2 n$ and $\mathbf{Q}(n)^{p, q}=0$ unless $p=q=-n$ in which case $\mathbf{Q}(n)^{-n,-n}=\mathbf{C}$. It is convenient to allow negative weights for this Hodge structure. The $n$-th Tate twist of $V$ is defined by $V(n):=V \otimes \mathbf{Q}(n)$, it is a Hodge structure of weight $k-2 n$ with $V(n)^{p, q}=V^{p+n, q+n}$.

Since the homomorphism $h_{n}$ acts via scalar multiplication on $V_{\mathbf{R}}$, it commutes with the representation $g_{\Sigma}^{m} h$ which defines $V_{m / 2}$. Therefore one has

$$
\left(V_{m / 2}\right)(n)=(V(n))_{m / 2} .
$$

This is also easy to see from the Hodge decompositions. Note that $V(-1) \neq\left(V_{-1 / 2}\right)_{-1 / 2}$ since the representations $h_{1}$ and $g_{\Sigma}^{2}$ are not equivalent.
2.7. The 'abstractly' defined half twists are in fact sub-Hodge structures of rather natural Hodge structures of CM-type. Recall that to the CM-field $K$ and the CM-type $\Sigma$ one can associate a weight one Hodge structure on the $\mathbf{Q}$-vector space $K$. These Hodge structures, and the abelian varieties associated to them, have been extensively investigated, cf. [La, DMOS]. If we endow $K$ with the trivial Hodge structure of weight zero, $K_{\mathbf{C}}=K^{0,0}$, this Hodge structure is just $K_{-1 / 2}$. One can identify $K_{-1 / 2}$ with $H^{1}\left(A_{K}, \mathbf{Q}\right)$ for an abelian variety $A_{K}$ with CM by the field $K$ and whose CM-type is the one used to define the (negative) half twist.

Let $V$ be a Hodge structure of CM-type with field $K$ of weight $n$, then $V \otimes_{\mathbf{Q}} K_{-1 / 2}$ is a Hodge structure of weight $n+1$ which has an action of the algebra $K \otimes_{\mathbf{Q}} K$. The element $x \otimes y \in K \otimes K$ acts as $(x \otimes y)(v \otimes z)=x v \otimes y z$ on $V \otimes_{\mathbf{Q}} K_{-1 / 2}$. We will identify some subspaces of $V \otimes_{\mathbf{Q}} K_{-1 / 2}$ with half twists of $V$.
2.8. Proposition. Let $(V, h, K)$ be a Hodge structure of CM-type. Then we have an inclusion of Hodge structures:

$$
V_{-1 / 2} \subset V \otimes_{\mathbf{Q}} K_{-1 / 2}
$$

more precisely:

$$
V_{-1 / 2}=\left\{w \in V \otimes_{\mathbf{Q}} K_{-1 / 2}:(x \otimes 1) w=(1 \otimes x) w, \quad \forall x \in K\right\}
$$

If $V$ admits a half twist, $V_{1 / 2}(-1)$ is also a sub-Hodge structure of $V \otimes_{\mathbf{Q}} K_{-1 / 2}$ :

$$
V_{1 / 2}(-1)=\left\{w \in V \otimes_{\mathbf{Q}} K_{-1 / 2}:(x \otimes 1) w=(1 \otimes \bar{x}) w, \quad \forall x \in K\right\}
$$

and the Hodge structure $V$ can be recovered from $V_{1 / 2}$ by applying the first result to $V_{1 / 2}$ rather then $V$ :

$$
V \subset V_{1 / 2} \otimes_{\mathbf{Q}} K_{-1 / 2}
$$

Proof. We give two proofs.
The field $K$ acts on the $\mathbf{Q}$-vector space $V$ and the (complex) eigenvalues of the action of $x \in K$ are the $\sigma_{j}(x) \in \mathbf{C}$, each with the same multiplicity. In particular, $x, \bar{x} \in K$ are eigenvalues of the $K$-linear extension of the action of $x$ on $V \otimes_{\mathbf{Q}} K$ and we denote by $V^{\prime}$ and $V^{\prime \prime}$ the corresponding eigenspaces:

$$
V \otimes_{\mathbf{Q}} K \cong V^{\prime} \oplus V^{\prime \prime} \oplus W, \quad \text { with } \quad(x \otimes 1) w= \begin{cases}(1 \otimes x) w & w \in V^{\prime} \\ (1 \otimes \bar{x}) w & w \in V^{\prime \prime}\end{cases}
$$

and $W$ is a $K$-stable complementary subspace. The projections on the summands give isomorphisms of $K=K \otimes 1$-vector spaces $V=V \otimes 1 \rightarrow V_{1}$ and $V \rightarrow V_{r+1}$. These $V_{j}$ are sub-Hodge structures of $V \otimes K_{-1 / 2}$ since the actions of $K$ and $h, g$ commute. The Hodge structure on $K_{-1 / 2}$ is defined by the $g(z) \in(K \otimes \mathbf{R})^{*}$ and thus $1 \otimes g(z)=g(z) \otimes 1$ on $V^{\prime}$ and $1 \otimes g(z)=g(\bar{z}) \otimes 1$ on $V^{\prime \prime}$. Note also that $g(t z)=t g(z)$ for $t \in \mathbf{R}$. Therefore

$$
h(z) \otimes g(z)=h(z) g(\bar{z}) \otimes 1=z \bar{z} h(z) g\left(z^{-1}\right) \otimes 1
$$

on $V^{\prime \prime}$ which identifies the Hodge structure on $V^{\prime \prime}$ with the one on $V_{1 / 2}(-1)$. The proof of $V^{\prime} \cong V_{-1 / 2}$ is similar but easier. The last result follows by applying the first result to $V_{1 / 2}$ rather than $V$ and using $\left(V_{1 / 2}\right)_{-1 / 2}=V$, note this identifies $V$ with a specific subspace of $V_{1 / 2} \otimes K_{-1 / 2}$.

Another proof of the first result is as follows: since for $x \in K$ the map $w \mapsto(1 \otimes x-x \otimes 1) w$ is a $\mathbf{Q}$-linear endomorphism of $V \otimes_{\mathbf{Q}} K_{-1 / 2}$, its kernel is a $\mathbf{Q}$-vector space and therefore

$$
V^{\prime}:=\left\{w \in V \otimes_{\mathbf{Q}} K_{-1 / 2}:(x \otimes 1) w=(1 \otimes x) w, \quad \forall x \in K\right\}
$$

is a $\mathbf{Q}$-vector space. Its complexification $V_{\mathbf{C}}^{\prime}$ is the folllowing subspace of $\left(V \otimes K_{-1 / 2}\right)_{\mathbf{C}}$ :

$$
V_{\mathbf{C}}^{\prime}=\oplus_{\sigma \in \Sigma \cup \bar{\Sigma}} V_{\sigma} \otimes K_{\sigma}
$$

with $V_{\sigma}$ (resp. $K_{\sigma}$ ) the subspace of $V_{\mathbf{C}}\left(\right.$ resp. $\left.K_{\mathbf{C}}\right)$ on which $x \in K$ acts as $\sigma(x)(\in \mathbf{C})$. Since $K_{-1 / 2}^{1,0}=\oplus_{\sigma \in \Sigma} K_{\sigma}$ we get:

$$
\left(V^{\prime}\right)^{r, s}=\left(\oplus_{\sigma \in \Sigma} V_{\sigma}^{r-1, s} \otimes K_{\sigma}\right) \oplus\left(\oplus_{\sigma \in \bar{\Sigma}} V_{\sigma}^{r, s-1} \otimes K_{\sigma}\right)
$$

which, since $K_{\sigma} \cong \mathbf{C}$, is just the definition of $V_{-1 / 2}^{r, s}$ hence $V^{\prime} \cong V_{-1 / 2}$. The other statements can be proved in a similar fashion.
2.9. Geometrical version of the half twist. The proposition shows that if $V$ is a sub-Hodge structure of $H^{k}(X, \mathbf{Q})$ for some projective variety $X$, then $V_{1 / 2}(-1)$ is a sub-Hodge structure of $H^{k}(X, \mathbf{Q}) \otimes H^{1}\left(A_{K}, \mathbf{Q}\right)$, which is itself a summand of $H^{k+1}\left(X \times A_{K}, \mathbf{Q}\right)$.
2.10. Polarizations. A polarization on the Hodge structure $(V, h)$ of weight $k$ is a bilinear map:

$$
\Psi: V \times V \longrightarrow \mathbf{Q}
$$

satisfying (for all $v, w \in V_{\mathbf{R}}$ ):

$$
\Psi(h(z) v, h(z) w)=(z \bar{z})^{k} \Psi(v, w)
$$

and

$$
\Psi(v, h(i) w) \text { is a symmetric and positive definite form: }
$$

$\Psi(v, h(i) w)=\Psi(w, h(i) v)$ for all $v, w \in V_{\mathbf{R}}$ and $\Psi(v, h(i) v)>0$ for all $v \in V_{\mathbf{R}}-\{0\}$. The primitive cohomology groups of algebraic varieties are polarized.

A polarized Hodge structure of CM-type with field $K$ is a polarized Hodge structure ( $V, h, \psi$ ) such that ( $V, h, K$ ) is of CM-type and such that

$$
\Psi(x v, w)=\Psi(v, \bar{x} w), \quad x \in K, v, w \in V
$$

2.11. A polarization on the half twist. The Hodge structure $K_{-1 / 2}$ has a polarization, cf. [La, DMOS. If $V$ has a polarization, then also $V \otimes K_{-1 / 2}$ has a polarization (the tensor product polarization) and by restriction one obtains a polarization on $V_{1 / 2}$.

One can also proceed more explicitly: if $\Psi$ is a polarization on $V$ then one chooses an element $\alpha \in K$ such that $\bar{\alpha}=-\alpha$ and such that for $\sigma \in \Sigma$ the purely imaginary complex numbers $\sigma(\alpha)$ all have positive imaginary part. Then the bilinear form

$$
\Psi^{\prime}: V \times V \longrightarrow \mathbf{Q}, \quad \Psi^{\prime}(v, w):=\Psi(v, \alpha w)
$$

is a polarization on $V_{1 / 2}$. To verify all the properties it is convenient to split

$$
\left(V_{1 / 2}\right)_{\mathbf{R}}=\oplus_{i=1}^{r}\left(V_{1 / 2}\right)_{i}
$$

corresponding to the decomposition $K \otimes_{\mathbf{Q}} \mathbf{R}=\oplus_{i=1}^{r} \mathbf{C}$.
2.12. Example. Let $F \in \mathbf{C}\left[X_{0}, \ldots, X_{n}\right]$, homogeneous of degree $n+2$, define a smooth variety $Y \subset \mathbf{P}^{n}$. Let

$$
X=Z \operatorname{eroes}\left(X_{n+1}^{n+2}-F\right) \quad\left(\subset \mathbf{P}^{n+1}\right), \quad \text { let } \quad V:=H^{n}(X, \mathbf{Q})_{0}
$$

be the primitive cohomology group of $X$. Then $V$ is a vector space over the field $K$ of $n+2$-th roots of unity, where the roots of unity act by multiplication on the variable $X_{n+1}$.

The vector space $H^{n, 0}(X)$ is one dimensional, hence we can apply the half twist to $V$ (for any CM -type which includes the embedding $\sigma$ of $K$ defined by $x v=\sigma(x) v$ for $v \in H^{n, 0}(X)$ ). Proposition 2.8 implies that

$$
V=H^{n}(X, \mathbf{Q})_{0} \hookrightarrow V_{1 / 2} \otimes K_{-1 / 2}
$$

with $V_{1 / 2}$ and $K_{-1 / 2}$ Hodge structures of weight $n-1$ and 1 respectively. In this case it is not hard to see geometrically that such Hodge structures exist.

Note that a general line $l \subset \mathbf{P}^{n}$ meets the hypersurface $Y:=\operatorname{Zeroes}(F)$ (the branch locus of the natural map $X \rightarrow \mathbf{P}^{n}$ ) in $n+2$ distinct points, so these have $n+2-3=n-1$ moduli. Since the grassmanian of lines in $\mathbf{P}^{n}$ has dimension $2(n-1)$, we find that for a general set of $n+2$ points on $\mathbf{P}^{1}$ there is a $n-1$-dimensional family $S$ of lines $l \subset \mathbf{P}^{n}$ each of which meet $Y$ in $n+2$ points with the same moduli. The union of these lines will be $\mathbf{P}^{n}$. Let $C_{n}$ be the
inverse image of (any) of these lines in $X$. After replacing $S$ by the desingularization of a finite cover $\tilde{S}$, we then get a dominant, rational map:

$$
\pi: \tilde{S} \times C_{n} \longrightarrow X
$$

Since a rational map is defined outside a subset of codimension at least 2, the pull-back of the regular $n$-form on $X$ extends to a regular $n$-form on $\tilde{S} \times C_{n}$. Hence we get $H^{n}(X, \mathbf{Q})_{0} \hookrightarrow$ $H^{n-1}(\bar{S}, \mathbf{Q}) \otimes H^{1}\left(C_{n}, \mathbf{Q}\right)$ where $\bar{S}$ is some compactified desingularization of $\tilde{S}$. We will relate this example to Shioda's results on Fermat type hypersurfaces in [VGI].

## 3. Kuga-Satake varieties

3.1. To a polarized Hodge structure $(V, \psi)$ of weight 2 with $\operatorname{dim} V^{2,0}=1$ the construction of Kuga and Satake associates a Hodge structure $\left(C^{+}(V), h_{s}\right)$ of weight 1 on the even Clifford algebra $C^{+}(V)$ of the quadratic space $(V, \psi)$ (see [KS and VG for a detailed construction). It has the property that there is an inclusion of Hodge structures $V \hookrightarrow C^{+}(V) \otimes C^{+}(V)$. The (isogeny class of) abelian variety associated to $C^{+}(V)$ is called the Kuga-Satake variety $K S(V)$ of $V$, so $H^{1}(K S(V), \mathbf{Q})=C^{+}(V)$.

In the remainder of this paper we consider such Hodge structures which are of CM-type with an imaginary quadratic field $K$. Since $\operatorname{dim} V^{2,0}=1$, the half twist $V_{1 / 2}$ is a Hodge structure and has weight one. Our main result is Theorem 3.10 which shows that $V_{1 / 2}$ is a summand of $\left(C^{+}(V), h_{s}\right)$. We also determine the other summands and relate them to half twists.

This completes the results of C. Voisin in [V], she already found that two summands of $C^{+}(V), S_{0}$ of dimension 2 and $S_{1}$ with $\operatorname{dim} S_{1}=\operatorname{dim} V$, such that $V \hookrightarrow S_{0} \otimes S_{1}$. We will identify $S_{1}$ with $V_{1 / 2}$ in Theorem 3.10. To find the simple summands of $\left(C^{+}(V), h_{s}\right)$ we use the Mumford-Tate group of the Hodge structure $V$.
3.2. The Mumford-Tate group. Recall that the Special Mumford-Tate group $S M T(V)$ of a polarized Hodge structure $(V, h, \psi)$ is the smallest algebraic subgroup $G$ of $G L(V)$, defined over $\mathbf{Q}$, for which $h(z) \in G(\mathbf{R})(\subset G L(V)(\mathbf{R}))$ for all $z \in \mathbf{C}$ with $z \bar{z}=1$.

The simple summands of the Hodge structure $\left(C^{+}(V), h_{s}\right)$ are then the irreducible subrepresentations of $S M T(V)$ in $C^{+}(V)$. It takes some rather long computations to determine these summands though.

The following lemma recalls the basic facts on the Mumford Tate group in this situation.
3.3. Lemma. Let $(V, h, \psi)$ be a weight 2 rational polarized Hodge structure of CM-type by an imaginary quadratic field $K=\mathbf{Q}(\phi)$, with $\bar{\phi}=-\phi$ and $\phi^{2}=-d$.
i) The Q-bilinear map:

$$
H: V \times V \longrightarrow K, \quad H(v, w):=\psi(v, w)+\phi^{-1} \psi(v, \phi w)
$$

is a hermitian form on the $K:=\mathbf{Q}(\phi)$-vector space $V($ so $\overline{H(v, w)}=H(w, v)$ and $H$ is $K$-linear in the second variable).
ii) We have

$$
S M T(V) \subset U(H):=\{g \in G L(V): H(g v, g w)=H(v, w), \quad g \phi=\phi g\}
$$

the unitary group of the $K$-vector space $V$ with the hermitian form $H$.
iii) In case ( $V, h$ ) has weight two, (the $\mathbf{R}$-linear extension of) the hermitian form $H$ has signature $\left(1 / 2 \operatorname{dim} V^{1,1}, \operatorname{dim} V^{2,0}\right)$ on the complex vector space $V \otimes_{\mathbf{Q}} \mathbf{R}$.

Proof. Since the weight is even, $\psi$ is symmetric and:

$$
\begin{aligned}
H(w, v) & =\psi(w, v)+\phi^{-1} \psi(w, \phi v) \\
& =\psi(v, w)+\phi^{-1} \psi(\phi v, w) \\
& =\psi(v, w)+(d \phi)^{-1} \psi\left(\phi^{2} v, \phi w\right) \\
& =\psi(v, w)-\phi^{-1} \psi(v, \phi w) \\
& =\frac{1(v, w)}{H( }
\end{aligned}
$$

The $K$-linearity in the second factor follows from:

$$
\begin{aligned}
H(v, \phi w) & =\psi(v, \phi w)+\phi^{-1} \psi\left(v, \phi^{2} w\right) \\
& =\psi(v, \phi w)-d \phi^{-1} \psi(v, w) \\
& =\phi\left(\psi(v, w)+\phi^{-1} \psi(v, \phi w)\right) \\
& =\phi H(v, w) .
\end{aligned}
$$

Using that $\psi(h(z) v, h(z) w)=(z \bar{z})^{k} \psi(v, w)$ where $k$ is the weight of $V$ and that $\phi h(z)=h(z) \phi$ we get:

$$
\begin{aligned}
H(h(z) v, h(z) w) & =\psi(h(z) v, h(z) w)+\phi^{-1} \psi(h(z) v, \phi h(z) w) \\
& =(z \bar{z})^{k} \psi(v, w)+\phi^{-1} \psi(h(z) v, h(z) \phi w) \\
& =(z \bar{z})^{k}\left(\psi(v, w)+\phi^{-1} \psi(v, \phi w)\right) \\
& =(z \bar{z})^{k} H(v, w),
\end{aligned}
$$

hence $S M T(V) \subset U(H)$.
In case the weight of $V$ is two, we write:

$$
V \otimes_{\mathbf{Q}} \mathbf{R}=V_{1} \oplus V_{2} \quad \text { with } \quad V_{1} \otimes_{\mathbf{R}} \mathbf{C}=V^{1,1}, \quad V_{2} \otimes_{\mathbf{R}} \mathbf{C}=V^{2,0} \oplus V^{0,2}
$$

Then $\psi$ is negative definite on $V_{2}$ and positive definite on $V_{1}$. Since $\phi$ and $h(z)$ commute, $\phi$ maps the eigenspaces $V^{p, q}$ into themselves hence $\phi\left(V_{i}\right) \subset V_{i}$. As $\phi^{2}=-d$ with $d>0, \phi$ does not have real eigenvalues and we can choose a $\mathbf{R}$-basis $f_{1}, \phi f_{1}, \ldots, f_{r}, \phi f_{r}$ of $V_{2}$ (and similarly for $\left.V_{1}\right)$. Since $\psi\left(f_{i}, \phi f_{i}\right)=d^{-1} \psi\left(\phi f_{i}, \phi^{2} f_{i}\right)=-\psi\left(\phi f_{i}, f_{i}\right)=-\psi\left(f_{i}, \phi f_{i}\right)$, we may assume this basis to be orthonormal.

Since $f_{1}, \ldots, f_{r}$ is a $\mathbf{C}=K \otimes_{\mathbf{Q}} \mathbf{R}$-basis of $V_{2}$ and $H\left(f_{i}, f_{j}\right)=\psi\left(f_{i}, f_{j}\right)$, we see that $H$ is negative definite on $V_{2}$ (and $H$ is positive definite on $V_{1}$ ).
3.4. Remark. If $(V, h, \psi)$ and $K=\mathbf{Q}(\phi)$ are as in section3.1, then the Lie group $U(H)(\mathbf{R})$ of real points of $U(H)$ is ismomorphic to $U(1, m-1)$.

For any $g \in U(H)(\mathbf{R})$ the Hodge structure $\left(V, h^{g}\right)$ with

$$
h^{g}: \mathbf{C}^{*} \longrightarrow G L(V)(\mathbf{R}), \quad z \longmapsto g h(z) g^{-1}
$$

is also polarized by the same $\psi$ and has $\phi \in \operatorname{End}_{H o d}\left(V, h^{g}\right)$. This implies that $S M T\left(V, h^{g}\right)=$ $U(H)$ for any general $g \in U(H)(\mathbf{R})$. One can show that the moduli space of such Hodge structures is isomorphic to the complex $m-1$-ball $U(1, m-1) /(U(1) \times U(m-1))$.

The inclusion $U(H) \subset S O(\psi)$ induces $U(1, m-1) \subset S O(2,2 m-2)$, this well-known inclusion is used for example to restrict modular forms from othogonal groups to unitary groups.
3.5. The universal cover of the orthogonal group $S O(\psi)$ has a natural spin representation on the even Clifford algebra $C^{+}(V)$. The following proposition describes the spin representation over $\mathbf{Q}$ for the quadratic forms under consideration. Over the complex numbers these results are very well known, but over a number field the situation is a bit delicate. The proposition will be used to decompose the spin representation as a representation of $U(H)$, the Mumford Tate group of a general Hodge structure of CM-type ( $V, h, \psi, K$ ) under consideration.
3.6. Proposition. Let $(V, h, \psi)$ be a polarized weight 2 Hodge structure with $\operatorname{dim} V^{2,0}=1$ which is of CM-type for an imaginary quadratic field $K$. Then there is a $\mathbf{Q}$-basis of $V$ such that

$$
\psi(x, x)=\sum_{i=1}^{m} d_{i} x_{i}^{2}+d \sum_{i=1}^{m} d_{i} x_{m+i}^{2}, \quad \text { with } \quad d_{1}<0, \quad d_{2}, \ldots, d_{m}>0
$$

The representation of so $(2 m)$ on $C^{+}(Q)$ decomposes as $C^{+}(V)=S^{2^{m-2}}$ where $S$ is a so $(\psi)$ representation of dimension $2^{m+1}$ whose irreducible components are:

$$
S \cong\left\{\begin{array}{llll}
S_{+} \oplus S_{-}, & \operatorname{End}_{s o(\psi}\left(S_{ \pm}\right) \cong D, & S_{ \pm} \otimes \mathbf{C} \cong \Gamma_{ \pm}^{2} & \text { if } m \equiv 0(4) \\
S_{1}^{2}, & \operatorname{End}_{s o(\psi)}\left(S_{1}\right) \cong K, & S_{1} \otimes \mathbf{C} \cong \Gamma_{+} \oplus \Gamma_{-} & \text {if } m \equiv 1,3(4)
\end{array}\right.
$$

with $D$ a skew field of degree 4 over $\mathbf{Q}$, and $\Gamma_{+}, \Gamma_{-}$are the two half-spin representations of so $(2 m) \otimes_{\mathbf{Q}} \mathbf{C}$, each of which has dimension $2^{m-1}$.

In case $m \equiv 2(4)$ there are two possibilities. If the equation $-\prod d_{i}=x^{2}+d y^{2}$ has a solution $(x, y) \in \mathbf{Q}^{2}$, then

$$
S \cong S_{+}^{2} \oplus S_{-}^{2}, \quad \operatorname{End}_{s o(\psi)}\left(S_{ \pm}\right) \cong \mathbf{Q}, \quad S_{ \pm} \otimes \mathbf{C} \cong \Gamma_{ \pm}
$$

(the split case). In case this equation has no solution we have

$$
S \cong S_{+} \oplus S_{-}, \quad E n d_{s o(\psi)}\left(S_{ \pm}\right) \cong D, \quad S_{ \pm} \otimes \mathbf{C} \cong \Gamma_{ \pm}^{2}
$$

(and $D$ is a skew field of degree 4 over $\mathbf{Q}$, we call this the non-split case).
Proof. To find this basis of $V$, choose $e_{1} \in V$ with $\psi\left(e_{1}, e_{1}\right) \neq 0$ and let $d_{1}:=\psi\left(e_{1}, e_{1}\right)$, $e_{m+1}:=\phi e_{1}$. Then $\psi\left(e_{m+1}, e_{m+1}\right)=d \psi\left(e_{1}, e_{1}\right)=d d_{1}$ and $d \psi\left(e_{1}, e_{m+1}\right)=\psi\left(\phi e_{1},(-d) e_{1}\right)=$ $-d \psi\left(e_{1}, e_{m+1}\right)$, hence $\psi\left(e_{1}, e_{m+1}\right)=0$. Next we take $e_{2} \in\left\langle e_{1}, e_{m+1}\right\rangle^{\perp}$ etc. Since the signature of $\psi$ is $(2-,(2 n-2)+)$, we may assume $d_{1}<0$ and $d_{2}, \ldots, d_{m}>0$.

We recall that over $K$ the spin representation of $s o(\psi)_{K}:=s o(\psi) \otimes_{\mathbf{Q}} K$ on $C^{+}(V)_{K}:=$ $C^{+}(V) \otimes_{\mathbf{Q}} K$ decomposes as a direct sum of $2^{m-1}$ copies of $\Gamma_{+} \oplus \Gamma_{-}$and then we take Galois invariants to find the irreducible so $(\psi)$-representations over $\mathbf{Q}$.

Let $V_{K}:=V \otimes_{\mathbf{Q}} \mathbf{Q}(\sqrt{-d})$, and consider the following $K$-basis of $V_{K}$ :

$$
f_{i}:=\frac{1}{2 d d_{i}}\left(\sqrt{-d} e_{i}+e_{m+i}\right), \quad f_{m+i}:=1 / 2\left(-\sqrt{-d} e_{i}+e_{m+i}\right), \quad(1 \leq i \leq m)
$$

where we wrote $\sqrt{-d} e_{i}$ for $e_{i} \otimes \sqrt{-d}$. One verifies that:

$$
\psi\left(\sum y_{j} f_{j}, \sum y_{j} f_{j}\right)=\sum_{i=1}^{m} y_{i} y_{m+i}
$$

Since $\psi\left(f_{j}, f_{j}\right)=0, \psi\left(f_{j}+f_{k}, f_{j}+f_{k}\right)=1$ if $|j-k|=m$ and is zero otherwise, we have in $C(V)_{K}$ :

$$
f_{j}^{2}=0, \quad f_{i} f_{m+i}+f_{m+i} f_{i}=1, \quad f_{j} f_{k}=-f_{k} f_{j} \quad \text { if }|j-k| \neq 0, m
$$

We denote conjugation on $K$ by a ${ }^{6,}$, this acts on $C(V)_{K}$ via the second factor and:

$$
\bar{f}_{i}=\frac{1}{d d_{i}} f_{m+i}, \quad \bar{f}_{m+i}=d_{i} d f_{i} .
$$

Let

$$
f:=f_{m+1} f_{m+2} \ldots f_{2 m}, \quad \text { then } \quad \bar{f}=\bar{f}_{m+1} \ldots \bar{f}_{2 m}=\left(d^{m} \prod_{i} d_{i}\right) f_{1} \ldots f_{m}
$$

In $C(V)_{K}$ we have:

$$
f \bar{f} f=\delta f, \quad \bar{f} f \bar{f}=\delta \bar{f}, \quad \text { with } \quad \delta:=(-1)^{m(m-1) / 2} d^{m} \prod_{i=1}^{m} d_{i},
$$

the second is just the conjugate of the first which is an easy computation:

$$
\begin{aligned}
f \bar{f} f & =\left(d^{m} \prod_{i} d_{i}\right)\left(f_{m+1} \ldots f_{2 m}\right)\left(f_{1} \ldots f_{m}\right)\left(f_{m+1} \ldots f_{2 m}\right) \\
& =\delta\left(f_{m+1} \ldots f_{2 m}\right)\left(f_{m} f_{m-1} \ldots f_{2} f_{1}\right)\left(f_{m+1} \ldots f_{2 m}\right) \\
& =\delta\left(f_{m+1} \ldots f_{2 m}\right)\left(f_{m} \ldots f_{2}\right)\left(1-f_{m+1} f_{1}\right)\left(f_{m+2} \ldots f_{2 m}\right) \\
& =\delta\left(f_{m+1} \ldots f_{2 m}\right)\left(f_{m} \ldots f_{3}\right)\left(1-f_{m+2} f_{2}\right)\left(f_{m+3} \ldots f_{2 m}\right)-0 \\
& =\ldots \\
& =\delta f .
\end{aligned}
$$

Now we consider the left $C(V)_{K}$-modules generated by $f$ and $\bar{f}$. These modules are isomorphic, in fact the relations we just proved imply that

$$
R_{\bar{f}}: C(V)_{K} f \longrightarrow C(V)_{K} \bar{f}, \quad x f \longmapsto x f \bar{f}
$$

is an isomorphism of left $C(V)_{K}$-modules with inverse

$$
R_{f}: C(V)_{K} \bar{f} \longrightarrow C(V)_{K} f, \quad y \bar{f} \longmapsto y \bar{f} f
$$

In [FH], Chapter 20, an inclusion $s o(\psi)_{K} \hookrightarrow C^{+}(V)_{K}\left(\subset C(V)_{K}\right)$ is constructed and it is shown (FH, 20.19 and 20.20) that, as an so $(\psi)_{K}$-module, $C(V)_{K} f$ is isomorphic to the direct sum of the two half-spin representations:

$$
C(V)_{K} f \cong \Gamma_{+, K} \oplus \Gamma_{-, K}, \quad \Gamma_{ \pm, K} \otimes_{K} \mathbf{C} \cong \Gamma_{ \pm} .
$$

Moreover, $C(V)_{K}^{+}$, the even Clifford algebra, is isomorphic to a product of two matrix algebras:

$$
C(V)_{K}^{+} \cong M_{2^{m-1}}(K) \times M_{2^{m-1}}(K)
$$

this implies that the spin representation of $s o(\psi)_{K}$ on $C^{+}(V)_{K}$ is isomorphic to $2^{m-1}$ copies of $\Gamma_{+, K} \oplus \Gamma_{-, K}$.

Before considering the situation over $\mathbf{Q}$ we recall that the center of $C^{+}(V)$ is $\mathbf{Q} \oplus \mathbf{Q} z$ with $z:=e_{1} e_{2} \ldots e_{2 m}$ and that $z^{2}=(-1)^{m} d^{m} \prod_{i} d_{i}^{2}(c f$. $\mathbb{[}, \S 5.2)$. Since $\sqrt{(-1)^{m} d^{m} \prod_{i} d_{i}^{2}} \in K$, the center of $C(V)_{K}^{+}$is $K \times K$ and one can verify that $\Gamma_{ \pm, K}$ are the two eigenspaces of $z$ in $C(V)_{K} f$. We also observe that $C(V)_{K} f \cap C(V)_{K} \bar{f}=\{0\}$, in fact if $a f=b \bar{f}$ then $a f \bar{f}=b \bar{f}^{2}=0$, hence $0=a f \bar{f} f=\delta a f$, hence $a f=0$.

The subspace $C(V)_{K} f$ is not defined over $\mathbf{Q}$ in general, but the direct sum

$$
S_{K}:=C(V)_{K} f \oplus C(V)_{K} \bar{f}
$$

obviously is defined over $\mathbf{Q}$, that is $S_{K}=S \otimes_{\mathbf{Q}} K$ for some $\mathbf{Q}$-vector space $S \subset C(V)$, in fact

$$
S=C(V)(f+\bar{f})+C(V) \sqrt{-d}(f-\bar{f}) .
$$

Moreover, $S$ is a representation space for $s o(\psi)\left(\subset \operatorname{so}(\psi)_{K}\right)$.

To decompose $S$ into irreducible components we determine $A:=\operatorname{End}_{s o(\psi)}(S)$, the endomorphisms of $S$ which commute with so $(\psi)$. Since $S_{K} \cong \Gamma_{+}^{2} \oplus \Gamma_{-}^{2}$, we have $A_{K}:=A \otimes_{\mathbf{Q}} K \cong$ $M_{2}(K) \times M_{2}(K)$, hence $\operatorname{dim}_{\mathbf{Q}} A=8$. It is clear that $A_{K}$ is generated by the center of $C^{+}(V)_{K}$ and the maps $R_{f}$ and $R_{\bar{f}}$. To determine $A$ it suffices to find the invariants under conjugation in $A_{K}$.

Obviously the center of $C^{+}(V)$ lies in $A$. Moreover, the maps

$$
\alpha, \beta: S \longrightarrow S, \quad \alpha: x \mapsto x(f+\bar{f}), \quad \beta: x \longmapsto x \sqrt{-d}(f-\bar{f})
$$

commute with $s o(\psi)$ (which acts from the left whereas $\alpha$ and $\beta$ act from the right). Note we have:

$$
(f+\bar{f})^{2}=f^{2}+f \bar{f}+\bar{f} f+\bar{f}^{2}=f \bar{f}+\bar{f} f=\delta
$$

the last equality holds since in $S_{K}$ we have:

$$
(a f+b \bar{f})(f \bar{f}+\bar{f} f)=0+a f \bar{f} f+b \bar{f} f \bar{f}+0=\delta(a f+b \bar{f})
$$

and similarly $(\sqrt{-d}(f-\bar{f}))^{2}=d(f \bar{f}+\bar{f} f)=d \delta$. Moreover, $\alpha \beta=-\beta \alpha$. Thus the $\mathbf{Q}$-algebra generated by $\alpha$ and $\beta$ is the quaternion algebra $D:=(\delta, d \delta)$ and

$$
A \cong D \otimes_{\mathbf{Q}} \mathbf{Q}(z), \quad \text { with } \quad D=(\delta, d \delta)
$$

Since $(\alpha \beta)^{2}=-\alpha^{2} \beta^{2}=-d \delta^{2}, D$ contains a copy of the field $\mathbf{Q}(\sqrt{-d}) \cong K$.
The center of $C^{+}(V)$ is $\mathbf{Q}(z)$ with $z^{2}=(-d)^{m} \prod_{i} d_{i}^{2}$, hence:

$$
\mathbf{Q}(z) \cong \mathbf{Q} \times \mathbf{Q} \quad \text { if } m \equiv 0(2), \quad \mathbf{Q}(z) \cong K \quad \text { if } m \equiv 1(2)
$$

As $d>0, d_{1}<0$ and $d_{2}, \ldots, d_{m}>0$, the sign of $\delta$ is the sign of $(-1)^{m(m-1) / 2}(-1)$. Thus $\delta, d \delta$ are both negative if $m \equiv 0,1$ (4) so $D \otimes_{\mathbf{Q}} \mathbf{R}$ is isomorphic to the algebra of quaternions, a skew field, and thus $D$ is also a skew field.

Hence in case $m \equiv 0$ (4) we have $A \cong D \times D$ with a skew field $D$, therefore $S$ splits up in two components $S_{+}$and $S_{-}$with $E n d_{s o(\psi)}\left(S_{ \pm}\right) \cong D$ and since $D \otimes_{\mathbf{Q}} \mathbf{C} \cong M_{2}(\mathbf{C})$ the $S_{ \pm} \otimes_{\mathbf{Q}} \mathbf{C}$ are both direct sums of two copies of one irreducible $s o(\psi)_{\mathbf{C}}$ representation.

In case $m \equiv 1(4), \mathbf{Q}(z) \cong K$ and $D$ contains a copy of $K$ hence $A \cong M_{2}(K)$ and thus $S$ is isomorphic to the sum of two isomorphic representations, irreducible over $\mathbf{Q}$, each of which, after tensoring with $K$, is the direct sum of two non-isomorphic representations.

In case $m \equiv 2(4)$, we have $A \cong D \times D$, but the structure of $D$ depends on the $d_{i}$. In fact, $D \cong(-d, \delta)$ (using $\alpha \beta$ and $\alpha$ as generators) and $\delta=-d_{1} d_{2} \ldots d_{m} d^{2 k}$ (with $2 k=m$ ), we also have $D \cong(-d, n)$ with $n:=-\prod_{i} d_{i} \in \mathbf{Z}_{>0}$. Hence $D \cong M_{2}(\mathbf{Q})$ iff $n=x^{2}+d y^{2}$ for some $x, y \in \mathbf{Q}^{2}$.

In case $m \equiv 3$ (4) we have, as in the case $m \equiv 1(4)$, that $A \cong M_{2}(K)$.
3.7. For the general $(V, h, \psi, K)$ we consider, the Mumford Tate group is $U(H)$. Thus the simple factors of the Hodge structure $\left(C^{+}(V), h_{s}\right)$ associated to $(V, h, \psi)$ are exactly the irreducible subrepresentations of the Lie algebra $u(H)(\subset s o(\psi))$ of $U(H)$ in $C^{+}(V)$ (which are defined over Q). We now determine the restriction of the $s o(\psi)$ representation $S$ from Proposition 3.6 to $u(H)$.
3.8. Proposition. Let $S$ be the $2^{m+1}$-dimensional so $(\psi)$-representation defined in Proposition 3.6. The $u(H)$-representation $S$ decomposes as follows:

$$
S \cong S_{0} \oplus S_{1} \oplus \ldots \oplus S_{m}, \quad S_{i} \cong S_{m-i}, \quad \operatorname{dim}_{\mathbf{Q}} S_{i}=2\binom{m}{i}
$$

The $S_{i}$ are irreducible $u(H)$-representations except $S_{l}$ if $2 l=m, l \equiv 2$ (4) and we are in the split case, in that case $S_{l} \cong\left(S_{l}^{\prime}\right)^{2}$ and $S_{l}^{\prime}$ is irreducible.

The $S_{i}$ are $K$-vector spaces and:

$$
\operatorname{End}_{u(H)}\left(S_{i}\right)= \begin{cases}K & (2 i \neq m) \\ D & (2 i=m)\end{cases}
$$

with $D$ a quaternion algebra (a skew field except for the split case) which contains $K$, but $\operatorname{End}_{u(H)}\left(S_{l}^{\prime}\right)=\mathbf{Q}$.

Proof. First we determine the inclusion $u(H)_{K} \hookrightarrow s o(\psi)_{K}$ and the restriction of the half-spin representations $\Gamma_{ \pm, K}$ to $u(H)_{K}$.

Extending the scalars from $\mathbf{Q}$ to $K$, the endomorphism $\phi$ of $V$ has two eigenspaces in $V \otimes_{\mathbf{Q}} K$ :

$$
V_{K}:=V \otimes_{\mathbf{Q}} K=V_{+} \oplus V_{-}, \quad(\phi \otimes 1) v= \pm(1 \otimes \phi) v \quad\left(v \in V_{ \pm}\right)
$$

Each of the eigenspaces is isotropic for the $K$-linear extension of $\psi$ to $V_{K}$, in fact,

$$
\begin{aligned}
\psi(v, w) & =d^{-1} \psi((\phi \otimes 1) v,(\phi \otimes 1) w) \\
& =d^{-1} \psi( \pm(1 \otimes \phi) v, \pm(1 \otimes \phi) w) \\
& =d^{-1} \phi^{2} \psi(v, w) \\
& =-\psi(v, w)
\end{aligned}
$$

The actions of $u(H)$ and $K$ on $V$ commute, hence $u(H)$ acts on the eigenspaces $V_{+}$and $V_{-}$, each of which has dimension $m$. In particular we have a Lie algebra map $u(H) \hookrightarrow g l\left(V_{+}\right)$, which, for dimension reasons, gives an isomorphism: $u(H)_{K} \cong g l\left(V_{+}\right)$. Since $\psi$ is preserved, this fixes the map $u(H) \rightarrow g l\left(V_{-}\right)$(in fact $\psi$ gives a duality $V_{-} \xlongequal{\cong} V_{+}^{*}$ ) and we get (with respect to the basis $f_{1}, \ldots, f_{2 m}$ of the proof of Prop. 3.6):

$$
u(H)_{K} \cong g l\left(V_{+}\right) \longrightarrow s o(\psi)_{K}, \quad A \longmapsto\left(\begin{array}{cc}
A & 0 \\
0 & -{ }^{t} A
\end{array}\right) .
$$

This inclusion $g l\left(V_{+}\right) \hookrightarrow s o(\psi)$ is the same as the one obtained from composing the isomorphism $g l\left(V_{+}\right) \cong \operatorname{end}\left(V_{+}\right) \cong V_{+}^{*} \otimes V_{-} \cong V_{-} \otimes V_{+}$and the inclusion (cf. [FH], formula (20.4)) $V_{-} \otimes V_{+} \hookrightarrow \wedge^{2} V_{K} \cong s o(\psi)$. ¿From [FH, 20.15 we obtain the restrictions of the half-spin representations $\Gamma_{ \pm, K}$ to $u(H)_{K}$ :

$$
\Gamma_{+, K} \cong \oplus_{i} \Lambda^{2 i} V_{+}, \quad \Gamma_{-, K} \cong \oplus_{i} \Lambda^{2 i+1} V_{+}
$$

(this is actually only an isomorphism of $s u(H)_{K}$-representations).
To find the irreducible representations of $s u(H)$ which are defined over $\mathbf{Q}$ we need to know how the conjugation on $K$ acts. For this we use the following $C(V)$-modules ( $\mathbb{F H}], 20.12$ ):

$$
C(V) f \cong \oplus_{i=0}^{m} \Lambda^{i} V_{+}, \quad C(V) \bar{f} \cong \oplus_{i=0}^{m} \Lambda^{i} V_{-}
$$

The $f_{i_{1}} \ldots f_{i_{r}} f$ with $1 \leq i_{1}<\ldots<i_{r} \leq m$ are a basis of $C(V)_{K} f$, and these correspond to the elements $f_{i_{1}} \wedge \ldots \wedge f_{i_{r}} \in \Lambda^{r} V_{+}$. Their conjugate is, up to a constant, $f_{i_{1}+m} \ldots f_{i_{r}+m} \bar{f}$ which
corresponds to $f_{i_{1}+m} \wedge \ldots \wedge f_{i_{r}+m} \in \Lambda^{r} V_{-} \cong \Lambda^{m-r} V_{+}$. In particular, we get conjugation-invariant subspaces

$$
S_{i, K}=\left(\Lambda^{i} V_{+}\right) f \oplus\left(\Lambda^{m-i} V_{+}\right) \bar{f} \quad\left(\subset S_{K}=C(V)_{K} f \oplus C(V)_{K} \bar{f}\right)
$$

for $0 \leq i \leq m$. Thus $S_{i, K}=S_{i} \otimes_{\mathbf{Q}} K$ for a subspace $S_{i} \in C(V)$ which is invariant under $s u(H)$. If $i \neq m-i$, the two summands are not isomorphic as $g l\left(V_{+}\right)_{K}$-representations, hence $S_{i}$ is an irreducible $s u(H)$-representation (if $W \subset S_{i}$ is an invariant subspace, then $W \otimes_{\mathbf{Q}} K$ is a $g l\left(V_{+}\right)_{K}$ and Galois invariant subspace of $S_{i, K}$ and hence $W=\{0\}$ or $\left.W=S_{i}\right)$. Since $\operatorname{End}_{s l\left(V_{+}\right)}\left(S_{i, k}\right)=K^{2}$ and $S_{i}$ is irreducible, $B:=\operatorname{End}_{s u(H)}\left(S_{i}\right)$ is a field, of degree two over Q. As $B \otimes_{\mathbf{Q}} K \cong K^{2}$ we get $B \cong K$. Note that $S_{i} \cong S_{m-i}$.

The interesting case is when $m$ is even and $2 l=m$. Both summands of $S_{l, K}$ are in the same half spin representation (hence in the same eigenspace of the center of $\left.C^{+}(V)_{K}\right)$. The maps $R_{f}$ and $R_{\bar{f}}$ generate $E n d_{s l\left(V_{+}\right)}\left(S_{l, K}\right) \cong M_{2}(K)$, hence $E n d_{s u(H)}\left(S_{l}\right)=(\delta, d \delta)$ (see the proof of Prop. 3.6). If $d \equiv 0(4)$ this quaternion algebra is a skew field and hence $S_{l}$ is irreducible. If $d \equiv 2(4), S_{l}$ is irreducible in the non-split case and in the split case is isomorphic to $\left(S_{l}^{\prime}\right)^{2}$ with $S_{l, K}^{\prime} \cong \Lambda^{l} V_{+}$. Thus we always have $\operatorname{End}_{s u(H)}\left(S_{l}\right) \cong D$, and $K \subset D$, but $D$ is not a skew field in the split case (then $D \cong M_{2}(\mathbf{Q})$ ). Since $S_{l, K}^{\prime}$ is irreducible we have $E n d_{s u(H)}\left(S_{l}^{\prime}\right) \cong \mathbf{Q}$.
3.9. The previous propositions show that the Kuga-Satake Hodge structure $\left(C^{+}(V), h_{s}\right)$ associated to ( $V, h, \psi, K$ ) decomposes as

$$
C^{+}(V) \cong\left(S_{0} \oplus S_{1} \oplus \ldots \oplus S_{m}\right)^{2^{m-2}}
$$

with $m=\operatorname{dim}_{K} V$ and they give the decomposition in simple factors as well as the endomorphism rings of the $S_{i}$ in the generic case.

The Hodge structure on the $S_{i}$ can be obtained as follows. Since $V$ is $K$-vector space, the exterior products $\wedge_{K}^{i} V$ are well-defined and, as $\mathbf{Q}$-vector spaces, they have dimension $2\binom{m}{i}$, which is just the dimension of the summand $S_{i}$. Weil already pointed out that there is a natural inclusion

$$
\wedge_{K}^{i} V \hookrightarrow \wedge^{i} V
$$

and the $\wedge_{K}^{i} V$ are sub-Hodge structures of $\wedge^{i} V$ of weight $2 i$. Combining Tate and half twists of these, one obtains weight 1 Hodge structures which are the summands of the Kuga-Satake Hodge structure. Note that the moduli of $V$, which has weight two, and of $S_{1}$ (with its $K$ action) which has weight one, are both the $m$-1-ball. Part of this theorem was already proved by Voisin in (V].
3.10. Theorem. Let $(V, h, \psi)$ be a polarized weight 2 Hodge structure with $\operatorname{dim} V^{2,0}=1$ which is of CM-type for an imaginary quadratic field $K$.

Then the Hodge structure on the summand $S_{i}$ (see Proposition 3.8) of the Kuga-Satake Hodge structure is:

$$
S_{i} \cong\left(\wedge_{K}^{i} V\right)(i-1)_{1 / 2}
$$

In particular, $S_{1}=V_{1 / 2}$ and $S_{0}$ is the CM-type Hodge structure of weight one on $K$. Moreover:

$$
V \hookrightarrow S_{0} \otimes S_{1}
$$

The dimensions of the eigenspaces for the $K$-action on $S_{i}^{1,0}$ are:

$$
\operatorname{dim} S_{i, \sigma}^{1,0}=\binom{m-1}{i-1}, \quad \operatorname{dim} S_{i, \bar{\sigma}}^{1,0}=\binom{m-1}{i}, \quad \text { with } \quad m=\operatorname{dim}_{K} V
$$

Proof. Let $V_{\mathbf{C}}=V_{+} \oplus V_{-}$be the decomposition in eigenspaces for the $K$-action, combining this with the Hodge decomposition we get:

$$
V_{\mathbf{C}}=V_{+}^{2,0} \oplus V_{+}^{1,1} \oplus V_{-}^{0,2} \oplus V_{-}^{1,1}
$$

We choose a basis $e_{1}, \ldots, e_{m}$ of $V_{+}$and $e_{m+1}, \ldots, e_{n}$ of $V_{-}$such that

$$
h(z)=\operatorname{diag}\left(z^{2}, 1 \ldots, 1, z^{-2}, 1, \ldots, 1\right) \quad(\in S O(\psi)(\mathbf{C})) \quad\left(z \in S^{1}\right)
$$

so, for example, $V^{0,2}=V_{-}^{0,2}=\left\langle e_{m+1}\right\rangle$. Then $h(z)$ lies in the 1-parameter subgroup generated by $H_{1}:=1 / 2\left(e_{1} \wedge e_{m+1}\right) \in \wedge^{2} V_{\mathbf{C}} \cong s o(\psi)_{\mathbf{C}}$. From the proof of FH], 20.15, one finds that $H_{1}$ multiplies $w:=e_{i_{1}} \wedge \ldots \wedge e_{i_{k}}, 1 \leq i_{1}<\ldots<i_{k} \leq m$, by $+1 / 2$ if $i_{1}=1$ and else by $-1 / 2$. Hence $h_{s}(z)$ multiplies $w$ by $z$ if $i_{1}=1$ and else by $\bar{z}$. Since:

$$
S_{i, \mathbf{C}}=\wedge^{i} V_{+} \oplus \wedge^{m-i} V_{+}=\left(V_{+}^{2,0} \otimes \wedge^{i-1} V_{+}^{1,1}\right) \oplus\left(\wedge^{i} V_{+}^{1,1}\right) \oplus\left(V_{+}^{2,0} \otimes \wedge^{m-i-1} V_{+}^{1,1}\right) \oplus\left(\wedge^{m-i} V_{+}^{1,1}\right)
$$

the action of $h_{s}(z)$ is by $\operatorname{diag}(z, \bar{z}, z, \bar{z})$. The isomorphism $\wedge^{m-i} V_{+} \cong \wedge^{i} V_{+}^{*} \cong \wedge^{i} V_{-}$induces

$$
V_{+}^{2,0} \otimes \wedge^{m-i-1} V_{+}^{1,1} \cong \wedge^{i} V_{-}^{1,1}, \quad \wedge^{m-i} V_{+}^{1,1} \cong V_{-}^{2,0} \otimes \wedge^{i-1} V_{-}^{1,1}
$$

Therefore the Hodge decomposition of $S_{i, \mathbf{C}} \cong \wedge^{i} V_{+} \oplus \wedge^{i} V_{-}$is given by:

$$
S_{i}^{1,0}=\left(V_{+}^{2,0} \otimes \wedge^{i-1} V_{+}^{1,1}\right) \oplus\left(\wedge^{i} V_{-}^{1,1}\right)
$$

The eigenspace decomposition of $S_{i}^{1,0}$ is now obvious and we see that $S_{1}=V_{1 / 2}$.
On the other hand, following Weil, we have

$$
\left(\wedge_{K}^{i} V\right) \otimes_{\mathbf{Q}} \mathbf{C}=\left(\wedge^{i} V_{+}\right) \oplus\left(\wedge^{i} V_{-}\right) \quad \hookrightarrow \quad \oplus_{j}\left(\wedge^{i-j} V_{+} \otimes \wedge^{j} V_{-}\right)=\left(\wedge^{i} V\right) \otimes \mathbf{C}
$$

The Hodge structure on $\wedge_{K}^{i} V$ induced from this inclusion is

$$
\left(\wedge_{K}^{i} V\right) \otimes_{\mathbf{Q}} \mathbf{C}=\left(V_{+}^{2,0} \otimes \wedge^{i-1} V_{+}^{1,1}\right) \oplus\left(\wedge^{i} V_{+}^{1,1} \oplus \wedge^{i} V_{-}^{1,1}\right) \oplus\left(\wedge^{i-1} V_{+}^{1,1} \otimes V_{-}^{0,2}\right)
$$

thus the Hodge numbers are $(2,0)+(i-1, i-1)=(i+1, i-1),(i, i)$ and $(i-1, i+1)$. Therefore if we Tate twist $(i-1)$-times and then do a half twist we obtain a Hodge structure of weight one which is just the one obtained from $h_{s}$.

The inclusion $V \subset V_{1 / 2} \otimes K(-1)_{1 / 2}=S_{1} \otimes S_{0}$ follows from Proposition 2.8.
3.11. Example. We construct, geometrically, a 9 dimensional family of polarized Hodge structures with $h^{2,0}=1, h^{1,1}=18$ with CM by the field $K \cong \mathbf{Q}(\sqrt{-3})$. For a Hodge structure $V$ of this family we identify the Hodge structures $S_{0}$ and $S_{1}$ as in Theorem 3.10 and we give a geometrical realisation of the inclusion $V \hookrightarrow S_{0} \otimes S_{1}$.

For $a_{1}, \ldots, a_{12} \in \mathbf{C}$ we define an (isotrivial) elliptic surface $S$ over $\mathbf{P}^{1}$ by the the Weierstrass model:

$$
S: \quad Y^{2}=X^{3}+\prod_{i=1}^{12}\left(t-a_{i}\right), \quad S \rightarrow \mathbf{P}^{1}, \quad(X, Y, t) \longmapsto t
$$

Since $S$ has twelve fibers which are cuspidal it is a K3 surface (and $\omega:=Y^{-1} \mathrm{~d} X \wedge \mathrm{~d} t$ is a no where zero holomorphic 2 -form on $S$ ).

The orthogonal complement in $H^{2}(S, \mathbf{Q})$ of the classes of a fiber and the section at infinity is a sub-Hodge structure $V$ of dimension 20 in with $V^{2,0}=1$ and the field $K=\mathbf{Q}(\sqrt{-3})$ acts on $V$ via the automorphism $(X, Y, t) \mapsto\left(\zeta^{2} X, Y, t\right)$ with a primitive 6 -th root of unity $\zeta$.

Define curves $C, C^{\prime}$, of genus 25 and 1 by:

$$
C: \quad y^{6}=\prod_{i=1}^{12}\left(x-a_{i}\right) ; \quad \quad C^{\prime}: \quad v^{2}=u^{3}-1
$$

Both of these curves have automorphisms of order 6:

$$
\psi: C \rightarrow C, \quad(x, y) \mapsto(x, \zeta y) ; \quad \psi^{\prime}: C^{\prime} \rightarrow C^{\prime}, \quad \psi^{\prime}:(u, v) \mapsto\left(\zeta^{2} u,-v\right)
$$

The surface $S$ is the (minimal model of the desingularisation of the) quotient of $C \times C^{\prime}$ by the automorphism $\phi=\left(\psi^{-1}, \psi^{\prime}\right)$ of order 6 , the quotient map is given by

$$
\pi: S=C \times C^{\prime} \longrightarrow S, \quad((x, y),(u, v)) \longmapsto(X, Y, t)=\left(y^{2} u, y^{3} v, x\right)
$$

To define $V_{1}$, consider the following rational 1-forms on $C$ :

$$
\omega_{a, b}:=x^{a} y^{b} \frac{\mathrm{~d} x}{y^{5}}, \quad \text { note } \quad \psi^{*}: \omega_{a, b} \mapsto \zeta^{b+1} \omega_{a, b}
$$

It is easy to check that the $\omega_{a, b}$ with $a, b \geq 0$ and $a+2 b \leq 8$ are a basis of $H^{0}\left(C, \omega_{C}\right)$. In particular, the eigenspace of $\psi^{*}$ with eigenvalue $\zeta$ has dimension 9 (and is spanned by the $\omega_{a, 0}$ with $0 \leq a \leq 8$ ) whereas the eigenspace with eigenvalue $\zeta^{-1}=\zeta^{5}$ has dimension 1 (and is spanned by $\omega_{0,4}$ ).

Let $V_{1} \subset H^{1}(C, \mathbf{Q})$ be the $\mathbf{Q}$-subspace on which the eigenvalues of $\zeta$ are primitive 6 -th roots of unity. Then $\operatorname{dim} V_{1}=20$, and the associated abelian variety is of Weil type $(1,9)$. Let $V_{0}:=H^{1}\left(C^{\prime}, \mathbf{Q}\right)$, note $\psi^{\prime}$ acts on $H^{0}\left(C^{\prime}, \omega_{C^{\prime}}\right)=\left\langle\omega^{\prime}:=\mathrm{d} u / v\right\rangle$ as $\zeta^{-1}$.

The pull-back $\pi^{*}$ maps $V$ into the $\phi$-invariants in $H^{1}(C, \mathbf{Q}) \otimes H^{1}\left(C^{\prime}, \mathbf{Q}\right)$ and it is easy to verify that these invariants are exactly the $\phi$-invariants in $V_{1} \otimes V_{0}$. For dimension reasons we then have:

$$
V \cong \pi^{*} V=\left(V_{0} \otimes_{\mathbf{Q}} V_{1}\right)^{\langle\phi\rangle} .
$$

Since $V_{0} \cong K_{-1 / 2}$, the half twist of this identity gives $V_{1 / 2} \subset K \otimes V_{1} \cong V_{1}^{\oplus 2}$, which implies that the half twist of $V$ is just $V_{1}$ :

$$
V_{1 / 2} \cong V_{1}
$$

and that $\pi^{*}$ is a geometrical realization of the Kuga-Satake correspondence.
The parameter space of 20 dimensional Hodge structures with CM by $K$ and $V^{2,0}=1$ is (a quotient of) the 9-ball (cf. 3.4). The K3 surfaces in this example are parametrized by 12 points in $\mathbf{P}^{1}$, Deligne and Mostov (DM) actually showed that the geometrical quotient $\left(\mathbf{P}^{1}\right)^{12} / / P G L(2)$ is a 9 -ball quotient. See [V] for old and new results on this moduli space.

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