Fourfolds of Weil type and the spinor map

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Abstract

Recent papers by Markman and O'Grady give, besides their main results on the Hodge conjecture and on hyperkähler varieties, surprising and explicit descriptions of families of abelian fourfolds of Weil type with trivial discriminant. They also provide a new perspective on the well-known fact that these abelian varieties are Kuga Satake varieties for certain weight two Hodge structures of rank six.

In this paper we give a pedestrian introduction to these results. The spinor map, which is defined using a half-spin representation of $SO(8)$, is used intensively. For simplicity, we use basic representation theory and we avoid the use of triality.
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Introduction

The recent papers [12,15] by Markman and O'Grady provide new descriptions of families of abelian fourfolds of Weil type. Markman uses these to prove that certain Hodge classes on these fourfolds are algebraic. Both show that these abelian varieties are isogeneous to the intermediate Jacobians of algebraic hyperkähler varieties of Kummer type. O'Grady further relates this to the Kuga Satake construction for the (primitive) second cohomology group of algebraic Kummer type varieties. See also [20] for further developments.

An abelian fourfold of Weil type has an imaginary quadratic field $K = \mathbb{Q}(\sqrt{-d})$ in its endomorphism algebra. Markman obtains the polarization on the abelian fourfold,

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an alternating form on H_1, from a symmetric(!) form on H_1 and the K-action. These fourfolds define two subspaces of the complexification of their first homology group H_1, a free $\mathbb{Z}$-module of rank 8. The first one is the $+i$-eigenspace of the complex structure on $H_1 \otimes \mathbb{R}$ defined by A. The second is one of the two eigenspaces of the K-action. These two subspaces turn out to be maximally isotropic subspaces for the symmetric form that determines the polarization.

In this paper we will mainly follow Markman's approach. He considers a free, rank 8, $\mathbb{Z}$-module V equipped with a symmetric bilinear form. This V will become the first cohomology group of the fourfolds of Weil type. The maximally isotropic subspaces of the complexification $V_{\mathbb{C}}$ of V are well-known to be parametrized by two copies of a Legendrian Grassmannian, a complex manifold of dimension six. The spinor map is a natural embedding of this Grassmannian in $\mathbf{P}^7$, the image is a quadric hypersurface Q^+. This map already made several appearances in algebraic geometry, for example in the study of vector bundles over hyperelliptic curves in [16], of K3 surfaces in [13], of secant varieties in [11] and of integrable systems [1].

The spinor map is best constructed using the representation theory of $Spin(V)$, a double cover of the orthogonal group $SO(V)$ defined by the bilinear form on V. The spin group has a half-spin representation whose projectivization is $\mathbf{P}^7$. The spinor map is equivariant for the action of $Spin(V)$. A natural integral structure on the half-spin representation allows one to identify it with the complexification of a free $\mathbb{Z}$ module S^+ of rank 8. There is a non-degenerate bilinear form on S^+ which defines the quadric Q^+.

A six dimensional analytic open subset $\Omega \subset Q^+$ parametrizes complex structures on $V_{\mathbb{R}}$ that preserve the bilinear form on V. Fixing a general element $s \in S^+$ and considering only the complex structures on $V_{\mathbb{R}}$ corresponding to $\ell \in \Omega \cap s^{\perp}$ produces a five dimensional family of complex tori $\mathcal{T}_{\ell}$, not algebraic in general, that have a Hodge class, called the Cayley class,

$$c_s \in H^{2,2}(\mathcal{T}_{\ell}, \mathbb{Z}) = H^4(\mathcal{T}_{\ell}, \mathbb{Z}) \cap H^{2,2}(\mathcal{T}_{\ell}).$$

The idea of using these tori and the associated action of $Spin(7) = Spin(s^{\perp})$ to study the Hodge conjecture for fourfolds of Weil type is due to V. Muñoz [14]. In Section 2.4 we observe that the existence of the Cayley classes can be deduced from a relation between the spinor and the Plücker map. Using representation theory we then compute the class c_s for certain s that are relevant for Markman's results in Proposition 3.4.

Triality, an automorphism of order three of $Spin(V)$, allows one to relate the standard representation of $Spin(V)$ (via $SO(V)$ on V) and the two half-spin representations, one of which is S^+. While it is prominent in [12], we use instead an 'ad hoc' Lemma 2.6. This lemma suffices to obtain the results on the Cayley class and for the Kuga Satake varieties.

For any $h \in S^+$ such that the sublattice $\langle h, s \rangle$ of S^+ spanned by h and s has rank two and is positive definite, the tori parametrized by the four dimensional domain $\Omega \cap \langle h, s \rangle^{\perp}$ turn out to be abelian fourfolds of Weil type. The imaginary quadratic field K depends on h and s, but fixing s and choosing h suitably, any such field occurs. The polarization is determined by K and the bilinear form on V. A further discrete invariant, the discriminant of a polarized abelian variety of Weil type, is always trivial for the fourfolds constructed in this way. See Theorem 4.6 for these results of Markman and O'Grady.

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The Hodge conjecture for an abelian fourfold A of Weil type is non-trivial. There is a natural 2-dimensional subspace $W_K \subset H^{2,2}(A, \mathbb{Q})$ of Hodge classes. It is not known in general if this subspace is spanned by classes of algebraic cycles. If c_s is algebraic, then all classes in W_K are also algebraic. Markman makes important progress in the study of the Hodge conjecture by showing that c_s is algebraic for all abelian fourfolds appearing in his construction, which are all fourfolds of Weil type with trivial discriminant. For this he uses deformation theory of sheaves on hyperkähler manifolds, see Section 5 for a brief outline.

We limit ourselves to a basic exposition of the constructions of Markman and O'Grady of the abelian fourfolds of Weil type with trivial discriminant and of the Cayley classes of Muñoz and Markman. The relation with the Kuga Satake construction is indicated in the last section.

0. Tori with an orthogonal structure

0.1. The lattice V

The complex tori we consider are all quotients of a fixed real vector space, with a varying complex structure, by a fixed lattice. Whereas one might expect an alternating form, a polarization, on the first cohomology group to be important, Markman instead fixes a symmetric, non-degenerate, bilinear form $(\bullet, \bullet)_V$ on a rank eight free $\mathbb{Z}$-module V of signature $(4+, 4-)$. He fixes a rank four free $\mathbb{Z}$-module W, defines $W^* := \mathrm{Hom}_{\mathbb{Z}}(W, \mathbb{Z})$ and

$$V := W \oplus W^*, \qquad \big((w_1, w_1^*), (w_2, w_2^*)\big)_V := w_1^*(w_2) + w_2^*(w_1).$$

If $e_1, \ldots, e_4$ is a $\mathbb{Z}$-basis of W and $e_{i+4} := e_i^*$, where $e_1^*, \ldots, e_4^*$ is the dual basis of W^* so that $e_i^*(e_j) = \delta_{ij}$ (Kronecker's delta), then

$$(v_1, v_2)_V := \sum_{i=1}^{4} x_i y_{i+4} + x_{i+4} y_i, \qquad \Big(v_1 := \sum_{i=1}^{8} x_i e_i, \quad v_2 := \sum_{i=1}^{8} y_i e_i \ \in V\Big),$$

hence $(V, (\bullet, \bullet)_V) \cong U^{\oplus 4}$, the direct sum of four copies of the hyperbolic plane $U = \big(\mathbb{Z}^2, \big(\begin{smallmatrix} 0 & 1 \\ 1 & 0 \end{smallmatrix}\big)\big)$.

In [12] and Section 5 one defines $W := H^1(X, \mathbb{Z})$ for an abelian surface X, but for the basic properties of the complex tori this is not needed.

0.2. Complex structures on $V_{\mathbb{R}}$

Let $V_{\mathbb{R}} := V \otimes_{\mathbb{Z}} \mathbb{R}$, it is an eight dimensional vector space over the real numbers. A complex structure on $V_{\mathbb{R}}$ is a linear map $J : V_{\mathbb{R}} \to V_{\mathbb{R}}$ with $J^2 = -I$. Such a map has two (complex) eigenspaces $Z_+, Z_- \subset V_{\mathbb{C}} := V \otimes_{\mathbb{Z}} \mathbb{C}$ corresponding to the eigenvalues $i, -i \in \mathbb{C}$ of J. These eigenspaces are complex conjugate, $\overline{Z_+} = Z_-$, where the complex conjugation on $V_{\mathbb{C}}$ is defined as $\overline{v \otimes z} = v \otimes \bar{z}$ for $v \in V$ and $z \in \mathbb{C}$.

$$V_{\mathbb{C}} = Z_+ \oplus Z_- = Z_+ \oplus \overline{Z_+}, \qquad J = (i, -i) \in End(Z_+) \oplus End(Z_-).$$

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Conversely, given two complex conjugate subspaces $Z_\pm \subset V_\mathbb{C}$ such that $V_\mathbb{C} = Z_+ \oplus Z_-$ one can define a linear map $\tilde{J} : V_\mathbb{C} \to V_\mathbb{C}$ by $\tilde{J}(v_+ + v_-) = iv_+ - iv_-$. Then there is a linear map $J : V_\mathbb{R} \to V_\mathbb{R}$ whose $\mathbb{C}$-linear extension to $V_\mathbb{C}$ is $\tilde{J}$. In fact, the inclusion $V_\mathbb{R} \hookrightarrow V_\mathbb{C}$ identifies $V_\mathbb{R}$ with the $(v_+, v_-) \in Z_+ \oplus Z_-$ with $\overline{v_+} = v_-$. Writing $v \in V_\mathbb{R}$ as $v = v_+ + \overline{v_+}$, one has $\tilde{J}v = iv_+ - i\overline{v_+} \in V_\mathbb{R}$, so J is just the restriction of $\tilde{J}$ to $V_\mathbb{R}$.

0.3. Orthogonal complex structures and isotropic subspaces

The $\mathbb{R}$-bilinear extension of $(\bullet, \bullet)_V$ defines a bilinear form on $V_\mathbb{R}$, denoted by the same symbol. We consider now the complex structures J that preserve this bilinear form, so $(Jv_1, Jv_2)_V = (v_1, v_2)_V$ for all $v_1, v_2 \in V_\mathbb{R}$. Equivalently, $J \in SO(V_\mathbb{R}, (\bullet, \bullet)_V)$ and we will call J an orthogonal complex structure. Notice that for such a complex structure J and for eigenvectors $v_{1+}, v_{2,+} \in Z_+$ we have, for the $\mathbb{C}$-bilinear extension of the bilinear form,

$$(v_{1+}, v_{2+})_V = (Jv_{1+}, Jv_{2+})_V = (iv_{1+}, iv_{2+})_V = i^2(v_{1+}, v_{2+})_V = -(v_{1+}, v_{2+})_V.$$

Hence the restriction of $(\bullet, \bullet)_V$ to Z_+ is identically zero. Thus Z_+ is an isotropic subspace of $V_\mathbb{C}$ (and since $\dim Z_+ = 4 = (1/2) \dim V_\mathbb{C}$ it is a maximally isotropic, or Legendrian, subspace of $V_\mathbb{C}$). Similarly Z_- is a maximally isotropic subspace of $V_\mathbb{C}$ (and since the bilinear form is non-degenerate it induces a duality $Z_+ \cong Z_-^*$).

One easily verifies that, conversely, an isotropic subspace $Z_+ \subset V_\mathbb{C}$ such that $V_\mathbb{C} = V_+ \oplus \overline{V_+}$ defines a complex structure J on $V_\mathbb{R}$ that preserves $(\bullet, \bullet)_V$. We summarize this in the following lemma.

0.4. Lemma

There is a natural bijection between the following two sets:

- the orthogonal complex structures $J \in SO(V_\mathbb{R}, (\bullet, \bullet)_V)$ on $V_\mathbb{R}$,
- the maximally isotropic subspaces Z of $V_\mathbb{C}$ such that $V_\mathbb{C} = Z \oplus \overline{Z}$.

1. The Legendrian Grassmannian and the spinor map

1.1.

We introduce a connected component $IG(4, V_\mathbb{C})^+$ of the Grassmannian of maximally isotropic subspaces of $V_\mathbb{C}$ in Section 1.7. It is isomorphic to a smooth six dimensional quadric $Q^+ \subset \mathbf{P}S_\mathbb{C}^+ \cong \mathbf{P}^7$, where $(S^+, (\bullet, \bullet)_{S^+})$ is a certain lattice of rank eight, see Section 1.10. This isomorphism is induced by the spinor map, which is equivariant for the action of the double cover $Spin(V)$ of $SO(V)$ on V and S^+ respectively. In Section 1.12 we identify a complex submanifold $\Omega \subset Q^+$ that parametrizes the orthogonal complex structures and we define complex tori $\mathcal{T}_\ell$ for $\ell \in \Omega$.

The spinor map was defined by Cartan [2] (see also [1]). The description given by Chevalley in [3] was used by Markman [11, Section 2]. We define the spinor map using the representation theory of orthogonal groups as in [5, Chapter 20] (but our $(v, w)_V$ is $2Q(v, w)$ in [5]).

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Recall that $V_{\mathbb{C}}$ is the complexification of V, a free $\mathbb{Z}$-module of rank $2n$ with $n = 4$, but in this section we will consider any positive integer n. The bilinear form $(\bullet, \bullet)_V$ is extended $\mathbb{C}$-bilinearly to $V_{\mathbb{C}}$ and the complexifications of $W, W^* \subset V$ are denoted by $W_{\mathbb{C}}, W_{\mathbb{C}}^*$. Whenever convenient we will also write $\mathbb{C}^{2n}$ for $V_{\mathbb{C}}$ and $SO(2n)$ for $SO(V_{\mathbb{C}})$ etc.

1.2. The Clifford algebra of $V_{\mathbb{C}}$

The Clifford algebra $C(V_{\mathbb{C}})$ of the complex vector space $V_{\mathbb{C}}$, of dimension $2n$, with the bilinear form $(\bullet, \bullet)_V$ is the quotient of the tensor algebra

$$C(V_{\mathbb{C}}) := \oplus_{k \geq 0} V_{\mathbb{C}}^{\otimes k} / \langle v \otimes w + w \otimes v - (v, w)_V \cdot 1 \rangle,$$

by the two sided ideal generated by the $v \otimes w + w \otimes v - (v, w)_V$ with $v, w \in V_{\mathbb{C}}$, or equivalently, by the two sided ideal generated by the $v \otimes v - (1/2)(v, v)_V$ for $v \in V$.

The Clifford algebra has dimension 2^{2n}. We identify $V_{\mathbb{C}}$ with its image in $C(V_{\mathbb{C}})$. The even Clifford algebra $C(V_{\mathbb{C}})^+$ is the image of $\oplus_{k \geq 0} V_{\mathbb{C}}^{\otimes 2k}$.

As in Section 0.1 we let

$$V_{\mathbb{C}} = W_{\mathbb{C}} \oplus W_{\mathbb{C}}^*, \qquad W_{\mathbb{C}} := \langle e_1, \ldots, e_n \rangle, \qquad W_{\mathbb{C}}^* := \langle e_{n+1}, \ldots, e_{2n} \rangle,$$

with $W_{\mathbb{C}}^* = \mathrm{Hom}(W_{\mathbb{C}}, \mathbb{C})$ the dual of $W_{\mathbb{C}}$, where $w^*(w) := (w, w^*)_V$ for $w \in W_{\mathbb{C}}, w^* \in W_{\mathbb{C}}^*$.

Since $W_{\mathbb{C}}, W_{\mathbb{C}}^*$ are isotropic one has $vw = -wv \in C(V_{\mathbb{C}})$ for all $v, w \in W_{\mathbb{C}}$ and also for all $v, w \in W_{\mathbb{C}}^*$. The subalgebras of $C(V_{\mathbb{C}})$ generated by $W_{\mathbb{C}}, W_{\mathbb{C}}^*$ are isomorphic to the exterior algebras $\wedge^\bullet W_{\mathbb{C}}$ and $\wedge^\bullet W_{\mathbb{C}}^*$ respectively.

Let $e^* := e_{n+1} \cdots e_{2n} \in C(V_{\mathbb{C}})$ be the product of the elements in a basis of $W_{\mathbb{C}}^*$. Then the left ideal $S_{\mathbb{C}} := C(V_{\mathbb{C}})e^*$ of $C(V_{\mathbb{C}})$ is isomorphic, as a $\mathbb{C}$ vector space, to $\wedge^\bullet W_{\mathbb{C}}$,

$$\sigma : \wedge^\bullet W_{\mathbb{C}} \xrightarrow{\cong} S_{\mathbb{C}} := C(V_{\mathbb{C}})e^*, \qquad w_1 \wedge w_2 \wedge \cdots \wedge w_r \longmapsto w_1 w_2 \ldots w_r e^*,$$

([3, II.2.2], [5, Exercise 20.12]). Under this isomorphism, left multiplication by $w \in W_{\mathbb{C}}$ and $w^* \in W_{\mathbb{C}}^*$ on $S_{\mathbb{C}}$ correspond to the following endomorphisms of $\wedge^\bullet W_{\mathbb{C}}$:

$$w\sigma(\eta) = \sigma(L_w \eta), \qquad w^* \sigma(\eta) = \sigma(D_{w^*}\eta), \qquad (\eta \in \wedge^\bullet W_{\mathbb{C}}),$$

where $L_w(\eta) := w \wedge \eta$ is left multiplication by w and D_{w^*} is the derivation on $\wedge^\bullet W_{\mathbb{C}}$ defined by

$$D_{w^*}(1) = 0, \qquad D_{w^*}(w_1 \wedge \cdots \wedge w_r) = \sum_{i=1}^{r} (-1)^{i-1} w^*(w_i)(w_1 \wedge \cdots \wedge \widehat{w_i} \wedge \cdots \wedge w_r),$$

(here $w^*(w) = (w, w^*)_V$ for $w \in W, w^* \in W_{\mathbb{C}}^*$).

These operations of $W_{\mathbb{C}}, W_{\mathbb{C}}^*$ on $\wedge^\bullet W_{\mathbb{C}}$ define a $C(V_{\mathbb{C}})$-module structure and σ is a homomorphism of $C(V_{\mathbb{C}})$-modules. It induces an isomorphism of $\mathbb{C}$-algebras between the even Clifford algebra and the direct sum of two matrix algebras (cf. [5, (20.13)])

$$C(V_{\mathbb{C}})^+ \cong \mathrm{End}(S_{\mathbb{C}}^+) \oplus \mathrm{End}(S_{\mathbb{C}}^-), \qquad S_{\mathbb{C}}^+ := \wedge^{even} W_{\mathbb{C}}, \qquad S_{\mathbb{C}}^- := \wedge^{odd} W_{\mathbb{C}}.$$

Since $\dim W_{\mathbb{C}} = n$ one has $\dim S_{\mathbb{C}}^\pm = 2^{n-1}$.

1.3. The spin group of $V_{\mathbb{C}}$

The conjugation on $C(V_{\mathbb{C}})$ is the anti-involution given by

$$x := x_1 \cdots x_r \longmapsto x^* := (-1)^r x_r \cdots x_1,$$

notice that it maps $C(V_{\mathbb{C}})^+$ into itself. The spin group of $V_{\mathbb{C}}$ is

$$Spin(V_{\mathbb{C}}) := \{x \in C(V_{\mathbb{C}})^+ : xx^* = 1, \quad x V_{\mathbb{C}} x^* \subset V_{\mathbb{C}}\}.$$

Elements in $Spin(V_{\mathbb{C}})$ thus induce linear maps on $V_{\mathbb{C}}$ and one has the following result.

1.4. Theorem

There is a surjective homomorphism of complex Lie groups

$$\rho_V : Spin(V_{\mathbb{C}}) \longrightarrow SO(V_{\mathbb{C}}), \qquad x \longmapsto [v \longmapsto x v x^*]$$

with kernel $\{\pm 1\}$.

Proof. For a proof see [5, Thm 20.28]. $\quad\square$

1.5. The half-spin representations

Besides this 'standard representation' of $Spin(V_{\mathbb{C}})$ on $V_{\mathbb{C}}$, one also has the two half-spin representations ρ^+, ρ^- of $Spin(V_{\mathbb{C}})$ on $S_{\mathbb{C}}^+$ and $S_{\mathbb{C}}^-$ respectively (vector spaces of dimension 2^{n-1}), given by left multiplication in $C(V_{\mathbb{C}})$:

$$\rho^\pm : Spin(V_{\mathbb{C}}) \longrightarrow GL(S_{\mathbb{C}}^\pm), \qquad x \longmapsto [\eta \longmapsto x\eta].$$

See [5, Exercise 20.38] for the fact that for $n \equiv 0 \mod 4$ the image of $Spin(V)$ lies in $SO(2^{n-1})$ (for a certain bilinear form β on $S_{\mathbb{C}}^+ \subset \wedge^\bullet W_{\mathbb{C}}$ also considered in [3, 3.2]). The center of $Spin(V_{\mathbb{C}})$, $\dim V_{\mathbb{C}} > 2$, is isomorphic to $(\mathbb{Z}/2\mathbb{Z})^2$ if n is even and is cyclic of order four otherwise (cf. [5, Exercise 20.36]). For n even, $n > 2$, the three quotients of $Spin(V_{\mathbb{C}})$ by the order two subgroups of the center are $SO(V_{\mathbb{C}})$ and the images of $Spin(V_{\mathbb{C}})$ in the two half-spin representations.

1.6. The Lie algebra $spin(V_{\mathbb{C}}) = so(2n)$

The Lie algebra $spin(V_{\mathbb{C}})$ of the subgroup $Spin(V_{\mathbb{C}})$ of the multiplicative group of $C(V_{\mathbb{C}})^+$ consists of the $x \in C(V_{\mathbb{C}})^+$ such that $x + x^* = 0$ and $xv - vx \in V_{\mathbb{C}}$ for all $v \in V_{\mathbb{C}}$ (cf. [3, p. 67–68]).

A basis of $spin(V_{\mathbb{C}})$ is given by the following $n(n-1)/2 + n(n-1)/2 + n^2 = n(2n-1)$ elements:

$$e_i e_j, \quad e_{i+n} e_{j+n} \quad \text{with} \quad 1 \le i \le j \le n; \qquad e_i e_{j+n} - \tfrac{1}{2} 1, \quad 1 \le i, j \le n.$$

To see that these elements are in $spin(V_{\mathbb{C}})$ (and to find their action on $V_{\mathbb{C}}$) one can use that for $x, y, v \in V_{\mathbb{C}}$ one has

$$[xy, v] := xyv - vxy = x(-vy + (y, v)_V) - (-xv + (x, v)_V)y = (y, v)_V x - (x, v)_V y.$$

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The Lie algebra $spin(V_{\mathbb{C}})$ is isomorphic to the Lie algebra $so(2n)$ of the orthogonal group $SO(V_{\mathbb{C}}) = SO(2n)$. This Lie algebra consists of the $X \in \operatorname{End}(V_{\mathbb{C}})$ such that $(Xv, w)_V + (v, Xw)_V = 0$ for all $v, w \in V_{\mathbb{C}}$. One finds that

$$so(2n) = \left\{ X = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \in \operatorname{End}(V_{\mathbb{C}}) : \quad A = -{}^t D, \quad {}^t B = -B, \quad {}^t C = -C \right\}.$$

An isomorphism $spin(V_{\mathbb{C}}) \to so(2n)$ is given by the differential of ρ_V, so by the representation of $spin(V_{\mathbb{C}})$ on $V_{\mathbb{C}}$ given by $x \cdot v := xv - vx$. Using the computation of $[xy, v]$ above and the notation $E_{i,j}$ for the $2n \times 2n$ elementary matrix whose only non-zero coefficient is $(E_{i,j})_{i,j} := 1$, one verifies that this isomorphism is given by

$$spin(V_{\mathbb{C}}) \xrightarrow{\cong} so(2n),$$
$$\begin{cases} e_i e_{n+j} & \longmapsto & X_{i,j}, & X_{i,j} & := & E_{i,j} - E_{n+j,n+i}, \\ e_i e_j & \longmapsto & Y_{i,j}, & Y_{i,j} & := & E_{i,n+j} - E_{j,n+i}, \\ e_{i+n} e_{j+n} & \longmapsto & Z_{i,j}, & Z_{i,j} & := & E_{n+i,j} - E_{n+j,i}. \end{cases}$$

We choose the Cartan subalgebra of $so(2n)$ to be the diagonal matrices in $so(2n)$ (as in [5, Section 18.1]):

$$\mathfrak{h} := \oplus_{i=1}^n \mathbb{C} H_i, \qquad H_i := E_{i,i} - E_{n+i,n+i}.$$

The dual $\mathfrak{h}^*$ of $\mathfrak{h}$ then consists of the linear maps (weights)

$$\mathfrak{h}^* := \oplus_{i=1}^n \mathbb{C} L_i, \qquad L_i \Big(\sum_{j=1}^n t_j H_j \Big) := t_i.$$

1.7. The isotropic Grassmannian

The (complex) n-dimensional subspaces of $V_{\mathbb{C}}$ are parametrized by the Grassmannian $G(n, V_{\mathbb{C}})$, which has dimension $n \cdot (2n - n) = n^2$. The maximally isotropic subspaces for $(\bullet, \bullet)_V$ (which are also those for the associated quadratic form) are parametrized by two (isomorphic, disjoint, connected) complex submanifolds of dimension $n(n - 1)/2$ of $G(n, V_{\mathbb{C}})$, denoted by $IG(n, V_{\mathbb{C}})^+$ and $IG(n, V_{\mathbb{C}})^-$. (See [6, Chapter 6] for linear subspaces of quadrics.) This generalizes the two rulings (families of lines) on a smooth quadric $Q \cong \mathbf{P}^1 \times \mathbf{P}^1$ in $\mathbf{P}^3$. We denote by $IG(n, V_{\mathbb{C}})^+$ the connected component which contains the maximally isotropic subspace $W_{\mathbb{C}}^*$. A complex maximally isotropic subspace Z defines a point $[Z] \in IG(n, V_{\mathbb{C}})^+$ if and only if $\dim(Z \cap W_{\mathbb{C}}^*) \equiv n \mod 2$. In case $n = 4$, the dimension of Z must be even and then also $[W_{\mathbb{C}}] \in IG(4, V_{\mathbb{C}})^+$.

We recall a local parametrization of $IG(n, V_{\mathbb{C}})^+$ by alternating $n \times n$ complex matrices. A basis of $W_{\mathbb{C}}^*$ is given by the last n basis vectors of V. Thus $W_{\mathbb{C}}^*$ is spanned by the columns of the $2n \times n$ matrix $\binom{0}{I}$. Slightly deforming $W_{\mathbb{C}}$, we obtain another subspace spanned by the columns of a $2n \times n$ matrix. Since $\det I = 1 \neq 0$ we may assume that the lower $n \times n$ submatrix is still invertible. Then we can find a basis of the same subspace given by the columns of a matrix $\binom{B}{I}$, the corresponding subspace will be denoted by Z_B. Thus we found a Zariski open subset $G(n, V_{\mathbb{C}})_0$ of $G(n, V_{\mathbb{C}})$ of dimension n^2 parametrized by $n \times n$ complex matrices.

In general Z_B will not be isotropic, but one easily verifies that

$$(\bullet, \bullet)_{V|Z_B \times Z_B} = 0 \quad \Longleftrightarrow \quad \begin{pmatrix} {}^t B & I \end{pmatrix} \begin{pmatrix} 0 & I \\ I & 0 \end{pmatrix} \begin{pmatrix} B \\ I \end{pmatrix} = 0 \quad \Longleftrightarrow \quad {}^t B + B = 0.$$

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Hence the vector space of alternating $n \times n$ matrices Alt_n provides us with a parametrization of a Zariski open subset of $IG(n, V_{\mathbb{C}})^+$ of dimension $n(n-1)/2$ which we denote by $IG(n, V_{\mathbb{C}})_0^+$:

$$Alt_n \xrightarrow{\cong} IG(n, V_{\mathbb{C}})_0^+ \hookrightarrow IG(n, V_{\mathbb{C}})^+, \qquad B \longmapsto [Z_B] = \left[\begin{pmatrix} B \\ I \end{pmatrix}\right].$$

The isotropic subspace Z_B is also the graph of the (alternating) map $W^* \to W$, $w^* \mapsto Bw^*$. Notice that $W^* = Z_0$.

The isotropic Grassmannian $IG(n, V_{\mathbb{C}})^+ = SO(V_{\mathbb{C}})/P$ is a homogeneous space where $P = P_{W^*}$ is the stabilizer of W^* in the group $SO(2n)$. The Lie algebra of P, which are the $X \in so(2n)$ with $XW^* \subset W^*$, consists of the $X \in so(2n)$ with $B = 0$.

1.8. The spinor map

We recall that the Pfaffian of an alternating $2m \times 2m$ matrix A is the complex number $\mathrm{Pfaff}(A)$ defined by the following identity in $\wedge^{2m}\mathbb{C}^{2m}$:

$$\mathrm{Pfaff}(A)e_1 \wedge \cdots \wedge e_{2m} = m!\omega_A^m, \qquad \left(\omega_A := \sum_{i<j} a_{ij}e_i \wedge e_j\right).$$

In the next theorem these Pfaffians appear as the coordinate functions of the spinor map $\gamma : IG(n, V_{\mathbb{C}})^+ \to \mathbf{P}S_{\mathbb{C}}^+$ that is moreover equivariant for the action of $Spin(V_{\mathbb{C}})$.

1.9. Theorem

Let $\rho^+ : Spin(V_{\mathbb{C}}) \to GL(S_{\mathbb{C}}^+)$ be the half-spin representation of $Spin(V_{\mathbb{C}})$ on $S_{\mathbb{C}}^+ = \wedge^{even} W_{\mathbb{C}} := \oplus_k \wedge^{2k} W_{\mathbb{C}}$.

(1) In case n is even, the highest weight α of $S_{\mathbb{C}}^+$ is $(L_1 + \cdots + L_n)/2$ and it is $(L_1 + \cdots + L_{n-1} - L_n)/2$ if n is odd.

(2) The one dimensional subspace

$$\langle 1 \rangle = \langle \wedge^0 W_{\mathbb{C}} \rangle \subset \wedge^{even} W_{\mathbb{C}}$$

is invariant under the Lie algebra of P. Thus there is a $Spin(V_{\mathbb{C}})$ equivariant map

$$\gamma : IG(n, V_{\mathbb{C}})^+ \longrightarrow \mathbf{P}S_{\mathbb{C}}^+, \qquad \gamma([\rho_V(\tilde{g})W_{\mathbb{C}}^*]) = \langle \rho^+(\tilde{g})1 \rangle$$

for $\tilde{g} \in Spin(V_{\mathbb{C}})$.

(3) For an alternating matrix $B \in M_n(\mathbb{C})$, let

$$X_B := \begin{pmatrix} 0 & B \\ 0 & 0 \end{pmatrix} \in so(2n), \qquad \tilde{g}_B := exp(X_B) \in Spin(V_{\mathbb{C}}).$$

In the standard representation $\rho_V : Spin(V_{\mathbb{C}}) \to SO(2n)$ one has

$$\rho_V(\tilde{g}_B) = \begin{pmatrix} I & B \\ 0 & I \end{pmatrix} \ (\in SO(2n)) \qquad \text{and} \quad \rho_V(\tilde{g}_B)Z_0 = Z_B.$$

In the half-spin representation on S^+ the action of $\tilde{g}_B$ is given by left multiplication:

$$\rho^+(\tilde{g}_B) : S_{\mathbb{C}}^+ \longrightarrow S_{\mathbb{C}}^+, \qquad \omega \longmapsto exp(\omega_B) \wedge \omega,$$

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and one has

$$exp(\omega_B) = \sum_{I, \sharp I \equiv 0 \mod 2} \mathrm{Pfaff}(B_I)e_I,$$

where I runs over the subsets of $\{1, \ldots, n\}$ with an even number of elements and $e_I = e_{i_1} \wedge \cdots \wedge e_{i_{2k}} \in \wedge^{even} W = S^+$ with $i_1 < \cdots < i_{2k}$.

(4) In the basis of $S_{\mathbb{C}}^+$ consisting of the e_I, the spinor map γ on the open subset $IG(n, V_{\mathbb{C}})_0^+$ is given by

$$\gamma : IG(n, V_{\mathbb{C}})_0^+ \longrightarrow \mathbf{P}S_{\mathbb{C}}^+, \qquad [Z_B] \longmapsto (\ldots : \mathrm{Pfaff}(B_I) : \ldots).$$

The image of γ is defined by quadrics.

Proof. The highest weight of the half-spin representation $S_{\mathbb{C}}^+$ is determined in [5, Proposition 20.15].

The Lie algebra of P is generated by the $X_{i,j}$ (matrices with $B = C = 0$) and the $Z_{i,j}$ (matrices with $A = B = D = 0$). The images of these elements in $End(S_{\mathbb{C}}^+)$ (as well as those of the $Y_{i,j}$) are:

$$so(2n) \longrightarrow End(S_{\mathbb{C}}^+),$$

$$\begin{cases} X_{i,j} & \longmapsto & e_i e_{n+j} - \frac{1}{2}\delta_{ij} & \longmapsto & L_{e_i} \circ D_{e_{n+j}} - \frac{1}{2}\delta_{ij}, \\ Y_{i,j} & \longmapsto & e_i e_j & \longmapsto & L_{e_i} \circ L_{e_j}, \\ Z_{i,j} & \longmapsto & e_{i+n} e_{j+n} & \longmapsto & D_{e_{i+n}} \circ D_{e_{j+n}}. \end{cases}$$

Since $D_{w^*}(1) = 0$ for all $w^* \in W^*$, we see that $X_{i,j}$ and $Z_{i,j}$ map 1 to an element in $\langle 1 \rangle$. Hence $Lie(P)$ maps $\langle 1 \rangle$ into itself and thus also the inverse image of P in $Spin(V_{\mathbb{C}})$ maps this line into itself.

The element $X_B \in so(2n)$ determined by B is $X_B = \sum_{i<j} b_{ij} Y_{i,j}$. It acts as left multiplication by $\omega_B := \sum b_{ij} e_i \wedge e_j$ on $\wedge^{even} W$ and thus $exp(X_B)$ is left multiplication by $exp(\omega_B) \in \wedge^{even} W$. The exponential map of an endomorphism α is $\sum \alpha^n/n!$. The multiplication between forms of even degree is commutative. The 2-form ω_B therefore generates a commutative subalgebra of nilpotent elements, hence $exp(X_B)$ is actually a finite sum. One also has

$$exp(\omega_B) = \prod_{i<j} exp(b_{ij} e_i \wedge e_j) = \prod_{i<j}(1 + b_{ij} e_i \wedge e_j).$$

We now show that, with B_I the submatrix of B with coefficients $b_{i,j}$ with $i, j \in I$,

$$exp(\omega_B) = \sum_{I, \sharp I \, even} \mathrm{Pfaff}(B_I)e_I.$$

In fact, $exp(\omega_B) \in \wedge^{even} W_{\mathbb{C}}$ is a linear combination of the $e_I = e_{i_1} \wedge \cdots \wedge e_{i_{2k}}$, where $i_1 \leq \cdots \leq e_{i_{2k}}$ $I = \{i_1, \ldots, i_{2k}\} \subset \{1, \ldots, n\}$ is a subset with an even number of elements. Since for an integer p one has that $\omega_B^p \in \wedge^{2p} W$, the coefficient of e_I is homogeneous of degree $2k$, with $2k = \sharp I$, in the coefficients b_{ij} of B and only those with $i, j \in I$ contribute. So the coefficient of e_I is determined by the $2k \times 2k$ alternating submatrix B_I of B with rows and columns indexed by I. Moreover this coefficient is $(\sum_{i_k < i_l, i_k, i_l \in I} b_{i_k i_l} e_{i_k} \wedge e_{i_l})^k/k!$, which is indeed $\mathrm{Pfaff}(B_I)$.

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Since $\rho_V(\tilde{g}_B)Z_0 = Z_B$ and $\gamma([Z_0]) = 1 \in S_\mathbb{C}^+$ we get $\gamma([Z_B]) = \langle \rho^+(\tilde{g})1 \rangle = \langle exp(\omega_B) \rangle \in S_\mathbb{C}^+ = \wedge^{even} W_\mathbb{C}$. The description of the spinor map follows immediately. For the equations defining the image see [3, III.3.2] or [9]. $\quad\square$

1.10. The spinor map for $n = 4$

In case $n = 4$, the spinor (or Cartan) map

$$\gamma : IG(4, V_\mathbb{C})^+ \longrightarrow \mathbf{P}S_\mathbb{C}^+$$

is an embedding whose image is a smooth quadric $Q^+ \subset \mathbf{P}S_\mathbb{C}^+$. We will often identify $IG(4, V_\mathbb{C})$ with Q^+ and simply write $[Z]$ for $\gamma([Z])$.

For Z_B in the open subset $IG(4, V_\mathbb{C})_0^+$, where $B = (b_{ij})$ is an alternating 4×4 matrix, this map is given, in a suitable basis of S^+, by (see Theorem 1.9.4):

$$\gamma : Z_B \longmapsto (z_1 : \ldots : z_8) =$$
$$\left(1 : b_{12} : b_{13} : b_{14} : b_{12}b_{34} - b_{13}b_{24} + b_{14}b_{23} : -b_{34} : b_{24} : -b_{23}\right).$$

The coordinate functions are, up to signs, the Pfaffians of the alternating submatrices of B with an even number of rows and columns. The closure of the image of γ is the spinor variety, a smooth quadric:

$$Q^+ = \gamma(IG(4, V_\mathbb{C})^+) = \{(z_1 : \ldots : z_8) \in \mathbf{P}S_\mathbb{C}^+ : \quad z_1z_5 + z_2z_6 + z_3z_7 + z_4z_8 = 0\}.$$

In fact the signs and the order of the coordinate functions on $S_\mathbb{C}^+$ were chosen in such a way as to obtain this simple equation.

The homogeneous coordinates above define a $\mathbb{Z}$-module $S^+ \cong \mathbb{Z}^8 \subset S_\mathbb{C}^+$ with bilinear form $(\bullet, \bullet)_{S^+}$ such that for $z = (z_1, \ldots, z_8) \in S^+$ one has $(z, z)_{S^+} = 2(z_1z_5 + z_2z_6 + z_3z_7 + z_4z_8)$. In particular, $S^+ \cong U^4$ and for $z \in S_\mathbb{C}^+$ one has $z \in Q^+$ iff $(z, z)_{S^+} = 0$ where we use the $\mathbb{C}$-bilinear extension of the bilinear form.

In this ad-hoc manner we obtain an integral structure on $S_\mathbb{C}^+$ and, as observed by Markman, one can define the spinor map over the integers.

1.11. Lemma

Let Z_1, Z_2 be two distinct maximally isotropic subspaces of $V_\mathbb{C}$ in the family parametrized by $IG(4, V_\mathbb{C})^+$. Then $Z_1 \cap Z_2 = \{0\}$ if and only if the complex line $\langle [Z_1], [Z_2] \rangle \subset \mathbf{P}S_\mathbb{C}^+$ is not contained in the spinor quadric Q^+.

Proof. Using the action of the orthogonal group, if $Z_1 \cap Z_2 = \{0\}$, then we can map Z_1, Z_2 to W, W^*. As $[W] = e_*, [W^*] = 1 \in S^+$ and $(e_*, 1)_{S^+} \neq 0$ it follows that the line $\langle [Z_1], [Z_2] \rangle$ is not contained in Q^+. On the other hand, if $Z_1 \cap Z_2 \neq 0$, then we may assume $Z_1 = W^*$ and $Z_1 = Z_B$ with B the rank two alternating 4×4 matrix with $\omega_B = e_1 \wedge e_2$. Then $[Z_1] = \langle 1 \rangle$ and $[Z_2] = \langle 1 + e_1 \wedge e_2 \rangle$ so that $\langle [Z_1], [Z_2] \rangle \subset Q^+$. $\quad\square$

1.12. Orthogonal complex structures and their period space Ω

We use the spinor map and the spinor variety Q^+ to parametrize the orthogonal complex structures and the complex tori that these define.

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An orthogonal complex structure J on $V_{\mathbb{R}}$ is determined by (and determines) a maximally isotropic subspace Z_+ such that $V_{\mathbb{C}} = Z_+ \oplus \overline{Z_+}$, see Lemma 0.4. Assume that $[Z_+] \in IG(4, V_{\mathbb{C}})^+$. Using the spinor map we see that $\ell := \gamma([Z_+])$ is a point of the quadric $Q^+ \subset \mathbf{P}S_{\mathbb{C}}^+$, that is $(\ell, \ell)_{S^+} = 0$. Since the spinor map is defined over $\mathbb{Q}$, we get $[\overline{Z_+}] = \bar{\ell}$, the complex conjugate of the point ℓ in $\mathbf{P}S_{\mathbb{C}}^+$. The condition that $Z_+ \cap \overline{Z_+} = 0$ is equivalent to the fact that the complex line spanned by $\ell, \bar{\ell}$ is not contained in Q^+ by Lemma 1.11. This again is equivalent to $(\ell, \bar{\ell})_{S^+} \neq 0$ and since $(\ell, \bar{\ell})_{S^+} \in \mathbb{R}$ we see that $(\ell, \bar{\ell})_{S^+}$ is either positive or negative.

We define an open (six dimensional, connected) analytic subset of Q^+ by

$$\Omega = \Omega_{S^+} := \{ \ell \in \mathbf{P}S_{\mathbb{C}}^+ : (\ell, \ell)_{S^+} = 0, \quad (\ell, \bar{\ell})_{S^+} > 0 \}.$$

Then any $\ell \in \Omega$ defines a maximal isotropic subspace Z_ℓ of $V_{\mathbb{C}}$ such that $V_{\mathbb{C}} = Z_\ell \oplus \overline{Z_\ell}$ and thus it defines an orthogonal complex structure J_ℓ on $V_{\mathbb{R}}$.

The complex structure J_ℓ on $V_{\mathbb{R}}$ defines a complex torus $\mathcal{T}_\ell$ of dimension four by requiring an isomorphism of weight 1 Hodge structures

$$H^1(\mathcal{T}_\ell, \mathbb{Z}) = (V, J_\ell), \qquad \text{i.e.} \quad H^{1,0}(\mathcal{T}_\ell) = Z_\ell.$$

This complex torus can also be defined as $\mathcal{T}_\ell = V_{\mathbb{C}}/(Z_\ell + V)$.

2. Tori with an orthogonal structure and a Cayley class

2.1.

In the previous section we used the spinor map to embed a Grassmannian $IG(4, V_{\mathbb{C}})^+$ of maximally isotropic subspaces in to $\mathbf{P}S_{\mathbb{C}}^+$. We now relate this embedding to the Plücker embedding of $Gr(4, V_{\mathbb{C}})$. For this we use again representation theory.

As a consequence, we find a natural map from S^+ to $\wedge^4 V$, the image of $s \in S^+$ is denoted by $c_s \in \wedge^4 V$. This is exploited as follows. For $\ell \in \Omega$ we defined a complex torus $\mathcal{T}_\ell$ and there is an isomorphism of Hodge structures $H^1(\mathcal{T}_\ell, \mathbb{Z}) = (V, J_\ell)$. Thus we can also identify the Hodge structures $\wedge^4 V = H^4(\mathcal{T}_\ell, \mathbb{Z})$. For $s \in S^+$ we then obtain a cohomology class $c_s \in H^4(\mathcal{T}_\ell, \mathbb{Z})$ which is Markman's Cayley class of s.

In Section 2.5 we recall Markman's result that the Cayley class is a Hodge class, so $c_s \in H^{2,2}(\mathcal{T}_\ell, \mathbb{Z})$, if and only if $\ell \in \Omega_{s^\perp} := s^\perp \cap \Omega$ where $s^\perp$ is the hyperplane in $S_{\mathbb{C}}^+$ defined by s using the bilinear form on S^+. Hence the five dimensional complex manifold $\Omega_{s^\perp}$ parametrizes the four dimensional complex tori with an orthogonal structure and Hodge class c_s.

2.2. The Plücker map

The Grassmannian $G(4, V_{\mathbb{C}})$ has a natural embedding, the Plücker map π, into a projective space $\mathbf{P}^N = \mathbf{P} \wedge^4 V_{\mathbb{C}}$ of dimension N, where $N + 1 = \binom{8}{4} = 70$:

$$\pi : G(4, V_{\mathbb{C}}) \longrightarrow \mathbf{P} \wedge^4 V_{\mathbb{C}}, \qquad Z \longmapsto [\wedge^4 Z].$$

The Plücker map is equivariant for the action of $GL(V_{\mathbb{C}})$.

On the open subset $G(4, V_{\mathbb{C}})_0$ of $G(4, V_{\mathbb{C}})$ defined in Section 1.7, the Plücker map is thus given by the determinants of the 4×4 submatrices of the 8×4 matrix $P := \binom{B}{I}$.

Using the basis of V from Section 0.1, the coefficient of $e_{i_1} \wedge \cdots \wedge e_{i_4}$ in

$$[\wedge^4 Z_B] = [r_1 \wedge \cdots \wedge r_4], \qquad \left(r_j = \sum_{k=1}^{8} P_{kj} e_k \in Z_B, \quad P := \binom{B}{I} \right)$$

is the determinant of the 4×4 submatrix of P with rows $i_1, \ldots, i_4$.

2.3. The spinor and the Plücker map

The theory of line bundles on homogeneous spaces provides a natural setting for the results below (cf. [5, Section 23.3, p.393], [1, Section II]), we only use basic representation theory. The Picard group of $G(4, V_{\mathbb{C}})$ is generated by the Plücker line bundle $\pi^* \mathcal{O}_{\mathbf{P}^N}(1)$. The restriction of this line bundle to $IG(4, V_{\mathbb{C}})^+$ does not generate the Picard group of $IG(4, V_{\mathbb{C}})^+$, but there is a line bundle $\mathcal{L}$ on $IG(4, V_{\mathbb{C}})^+$ such that

$$(\pi^* \mathcal{O}_{\mathbf{P}^N}(1))_{|IG(4, V_{\mathbb{C}})^+} \cong \mathcal{L}^{\otimes 2},$$

and $Pic(IG(4, V_{\mathbb{C}})^+) \cong \mathbb{Z}$ is generated by $\mathcal{L}$. One has $H^0(IG(4, V_{\mathbb{C}})^+, \mathcal{L}) = S_{\mathbb{C}}^+$ and the spinor map γ is the map defined by the global sections of $\mathcal{L}$.

From the isomorphism $\pi^* \mathcal{O}_{\mathbf{P}^N} \cong \mathcal{L}^{\otimes 2}$ over $IG(4, V_{\mathbb{C}})^+$, one can deduce that the Plücker map on $IG(4, V_{\mathbb{C}})^+$ is the composition of the spinor map γ with the second Veronese map ν on $\mathbf{P}S_{\mathbb{C}}^+$. The Veronese map is induced by

$$\nu : S^+ \longrightarrow Sym^2(S^+), \qquad s \longmapsto s \odot s.$$

More precisely, the group $Spin(V_{\mathbb{C}})$, a double cover of $SO(V_{\mathbb{C}})$, has a natural (half-spin) representation ρ^+ on $S_{\mathbb{C}}^+$ and on the 36-dimensional vector space $Sym^2(S_{\mathbb{C}}^+)$. This latter representation is reducible, due to the $Spin(V)$-invariant quadric on S^+ which dually defines an invariant one dimensional subspace Γ_0 of $Sym^2(S_{\mathbb{C}}^+)$. A complement of this subspace turns out to be an irreducible $Spin(V_{\mathbb{C}})$-representation, of dimension 35, and is denoted by $\Gamma_{2\alpha}$ ([3, Section 3.4], [5, Exercise 19.6]):

$$Sym^2(S_{\mathbb{C}}^+) \cong \Gamma_{2\alpha} \oplus \Gamma_0.$$

The subspace $\Gamma_{2\alpha}$ is spanned by the symmetric tensors $z \odot z \in Sym^2(S_{\mathbb{C}}^+)$ with $[z] \in Q^+ \subset \mathbf{P}S_{\mathbb{C}}^+$.

There is a decomposition of the 70-dimensional $\wedge^4 V_{\mathbb{C}}$ in two irreducible $Spin(V_{\mathbb{C}})$-representations of dimension 35 (it corresponds to the decomposition of $\wedge^4 V_{\mathbb{C}}$ into dual and anti-selfdual 4-forms for the Hodge star operator defined by $(\bullet, \bullet)_V$) ([5, Remarks(ii), p. 289–290]):

$$\wedge^4 V_{\mathbb{C}} = \Gamma_{2\alpha} \oplus \Gamma_{2\beta}.$$

The image of Q^+ by the second Veronese map spans the linear subspace $\mathbf{P}\Gamma_{2\alpha} \subset \mathbf{P}^N = \mathbf{P} \wedge^4 V_{\mathbb{C}}$.

Since on the open subset of $IG(n, 2n)^+$ parametrized by alternating matrices the spinor map is given by Pfaffians and the Plücker map is given by minors, this result implies that any quadratic expression in Pfaffians is a linear combination of minors, see [1].

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2.4. The Cayley classes

A remarkable consequence of the relation between the $Spin(V)$-representations $Sym^2(S^+)$ and $\wedge^4 V$ is that any element $s \in S^+$ defines a 4-form $c_s \in \wedge^4 V$, which is called the Cayley class of s ([12, Remark 12.4], [14, Section 2.1]). It is obtained as the composition

$$S^+ \xrightarrow{\ v\ } Sym^2(S^+) \cong \Gamma_{2\alpha} \oplus \Gamma_0 \longrightarrow \Gamma_{2\alpha} \longrightarrow \wedge^4 V, \qquad s \longmapsto c_s.$$

This map is equivariant for the action of $Spin(V)$. The stabilizers in $Spin(V)$ of s and c_s thus have the same Lie subalgebra. If $(s, s)_{S_+} \neq 0$, the complexification of this Lie algebra is isomorphic to $so(7)_{\mathbb{C}}$.

2.5. The Cayley class and Hodge classes

Let $\ell \in \Omega \subset Q^+$ and let $\mathcal{T}_\ell$ be the associated complex torus. The Hodge decomposition of the first cohomology group $H^1(\mathcal{T}_\ell, \mathbb{Z}) = (V, J_\ell)$ is given by the eigenspaces $Z_\ell, \overline{Z}_\ell = Z_{\bar{\ell}}$ of the orthogonal complex structure J_ℓ in $V_{\mathbb{C}}$:

$$H^1(\mathcal{T}_\ell, \mathbb{C}) = V_{\mathbb{C}} = Z_\ell \oplus Z_{\bar{\ell}}, \qquad J_\ell = (i, -i) \in End(Z_l) \oplus End(Z_{\bar{\ell}}).$$

To describe the Hodge structure on $H^k(\mathcal{T}_\ell, \mathbb{Z})$ we use the homomorphism

$$h_{V,\ell} : \ U(1) := \{z \in \mathbb{C} : z\bar{z} = 1\} \longrightarrow GL(V_{\mathbb{R}}), \qquad h_{V,\ell}(a + bi) := aI + bJ_\ell,$$

where $a, b \in \mathbb{R}$, $a^2 + b^2 = 1$. Notice that $aI + bJ_\ell = (a+bi, a-bi) \in End(Z_l) \oplus End(Z_{\bar{\ell}})$.

Since $H^k(\mathcal{T}_\ell, \mathbb{Z}) = \wedge^k H^1(\mathcal{T}_\ell, \mathbb{Z}) = \wedge^k V$, the Hodge decomposition $H^k(\mathcal{T}_\ell, \mathbb{C}) = \oplus H^{p,q}(\mathcal{T}_\ell)$ is defined by

$$H^{p,q}(\mathcal{T}_\ell) = \left(\wedge^p Z_\ell \right) \otimes \left(\wedge^q Z_{\bar{\ell}} \right) = \{x \in \wedge^k V_{\mathbb{C}} : \ h_{V,\ell}(a + bi) \cdot x$$
$$= (a + bi)^p (a - bi)^q x \quad \forall a + bi \in U(1)\},$$

In particular, the Hodge classes in $H^{2p}(\mathcal{T}_\ell, \mathbb{Z})$ are the invariants of the one-parameter subgroup $h_{V,\ell}$ of $SO(V_{\mathbb{R}})$.

The homomorphisms $h_{V,\ell}$ can be lifted to $Spin(V_{\mathbb{C}})$ and the following lemma describes the action of such a lift on $S_{\mathbb{C}}^+$. It implies that the weight one Hodge structure (V, J_ℓ) defines a weight two Hodge structure on S^+. A rank six Hodge substructure $H = H_\ell \subset S^+$ will be studied in Section 6.1.

It should be noted that if $\dim V_{\mathbb{C}} \neq 8$ then $Spin(V_{\mathbb{C}})$ only allows one to relate polarized weight two Hodge structures on V to complex structures on S^+ and on the even Clifford algebra $C(V)^+$. The special feature in the case $\dim V = 8$ is triality, an automorphism of order three of $Spin(V)$, which allows one to permute the three irreducible 8-dimensional representations V, S^+, S^-, see [5, Section 20.3], [3, Chapter 4], and which is implicit in the proof of Lemma 2.6.

2.6. Lemma

Let $V = U \oplus U^*$ be a decomposition of $V = \mathbb{C}^8$ with two maximally isotropic subspaces with $[U], [U^*] \in IG(4, V_{\mathbb{C}})^+$. For $t \in \mathbb{C}$, $t \neq 0$, the orthogonal transformation

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$(t\,\mathrm{id}_U, t^{-1}\mathrm{id}_{U^*}) \in (\mathrm{End}(U) \oplus \mathrm{End}(U^*)) \cap SO(V)$ has a lift $h(t) \in Spin(V)$ which acts as follows on S^+:

$$\rho^+(h(t))\,\ell_U = t^2\ell_U, \quad \rho^+(h(t))\,\ell_{U^*} = t^2\ell_{U^*}, \qquad \rho^+(h(t))\,s = s,$$
$$\forall s \in \langle \ell_U,\, \ell_{U^*} \rangle^\perp,$$

where $\ell_U, \ell_{U^*} \in S^+$ are (any) representatives of $\gamma([U]), \gamma([U^*]) \in \mathbf{P}S^+$.

Proof. We use that the spinor map is equivariant for the action of $Spin(V)$. There is an element of $Spin(V)$ mapping U to W since $IG(4,8)^+ = SO(V)/P$. Then U^* is mapped to Z_B for some $B \in Alt_4$ and it is easy to see that there is another element in $Spin(V)$ fixing W (so with $C = 0$) and mapping Z_B to $Z_0 = W^*$. We thus may replace W, W^* with U, U^*. The one parameter subgroup h acts as multiplication by t on $U \subset V$, hence h is generated by an $X \in \mathfrak{h} \subset spin(V)$ with $L_i(X) = 1$ for $i = 1, \ldots, 4$ (and thus $X = \sum H_i$). The weights of S^+ are $(\pm L_1 \pm L_2 \pm L_3 \pm L_4)/2$ with an even number of - signs, hence their values on X are $2, -2$, with multiplicity one, and 0 with multiplicity six. Thus $\rho^+(h(t))$ is semisimple with eigenvalues t^2, t^{-2} and 1, the last with multiplicity six. The eigenvalue t^{-2}, the lowest weight of S^+, is on Z_{U^*}, see Theorem 1.9. The element $g \in SO(V)$ that maps $e_i \mapsto e_{i+4}, e_{i+4} \mapsto e_i$ for $i = 1, \ldots, 4$ interchanges U and U^* and acts (in the Adjoint representation) as $-id$ on $\mathfrak{h}$, hence the eigenvalue t^2 must be on Z_U. As $Spin(V)$ preserves $(\bullet, \bullet)_{S^+}$, the decomposition into these eigenspaces is orthogonal. (For any n, the one parameter subgroup of $SO(V)$ that acts as multiplication by t^2, t^{-2} on e_1, e_{n+1} respectively and is the identity on $\langle e_1, e_{n+1} \rangle^\perp$ is generated by an $X \in spin(V)$ with $L_1(X) = 2, L_i(X) = 0$ for $i \geq 2$ and thus $(1/2)(\pm L_1 \pm L_2 \ldots \pm L_n)(X) = \pm 1$, showing that the lift of this subgroup to $Spin(V)$ has only eigenvalues t, t^{-1} on S^+, with the same multiplicities, and the same holds for S^-. A similar result holds for $SO(V)$ and its spin representation if $\dim V = 2n + 1$.) $\square$

We use this lemma to identify the complex tori $\mathcal{T}_\ell$ for which the Cayley class c_s is of Hodge type $(2,2)$. The following proposition is essentially [12, Lemma 12.2].

2.7. Proposition

Let $c_s \in \wedge^4 V$ be the Cayley class defined by $s \in S^+$, the integral lattice, and let $\ell \in \Omega_{S^+}$. Then c_s is an integral Hodge class in $H^4(\mathcal{T}_\ell, \mathbb{Z})$ exactly when $(\ell, s)_{S^+} = 0$:

$$c_s \in H^{2,2}(\mathcal{T}_\ell, \mathbb{Z}) := H^4(\mathcal{T}_\ell, \mathbb{Z}) \cap H^{2,2}(\mathcal{T}_\ell) \qquad \text{if and only if}$$
$$\ell \in \Omega_{s\perp} := \{\ell \in \Omega : (\ell, s)_{S^+} = 0\}.$$

Proof. First we observe that $h_{V,\ell}(z) \in SO(V_\mathbb{R})$ for all $z \in U(1)$. In fact, for $v, w \in V_\mathbb{R}$ we have

$$\big((aI + bJ_\ell)v, (aI + bJ_\ell)w\big)_V = a^2(v, w)_V + ab\big((v, J_\ell w)_V + (J_\ell v, w)_V\big)$$
$$+ b^2\big(J_\ell v, J_\ell w\big)_V = (v, w)_V,$$

because $(J_\ell v, J_\ell w)_V = (v, w)_V$ implies $(v, J_\ell w) = (J_\ell v, J_\ell^2 w)_V$ and $J_\ell^2 = -I$.

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The homomorphism lifting the one-parameter subgroup $h_{V,\ell} : U(1) \to SO(V_{\mathbb{C}})$ to $Spin(V_{\mathbb{C}})$ is denoted by

$$h_\ell : U(1) := \{z \in \mathbb{C} : z\bar{z} = 1\} \longrightarrow Spin(V_{\mathbb{C}}).$$

The action of $h_\ell(z) \in Spin(V_{\mathbb{C}})$ in the half-spin representation ρ^+ on $S_{\mathbb{C}}^+$ is (see Lemma 2.6):

$$\rho^+(h_\ell(z))\ell = z^2\ell, \quad \rho^+(h_\ell(z))\bar{\ell} = \bar{z}^2\bar{\ell}, \qquad \rho^+(h_\ell(z))s = s, \quad \forall s \in \langle \ell, \bar{\ell} \rangle^\perp.$$

Using the induced action of $Spin(V_{\mathbb{C}})$ on $s \odot s \in Sym^2(S_{\mathbb{C}}^+)$ and its image $c_s \in \wedge^4 V_{\mathbb{C}} = H^4(\mathcal{T}_\ell, \mathbb{C})$ we see that c_s is invariant under $h_\ell(z)$ for all $z \in U(1)$ if and only if s is invariant, so $s \in \langle \ell, \bar{\ell} \rangle^\perp$. For $s \in S^+$ the condition $(s, \ell)_{S+} = 0$ implies, by complex conjugation, that also $(s, \bar{\ell})_{S+} = 0$, which proves the proposition. $\square$

3. The Cayley class as $Spin(7)$-invariant

3.1.

In Section 2.4 we defined the Cayley class $c_s \in \wedge^4 V$ for $s \in S^+$. We compute this class explicitly in Proposition 3.4 for certain $s \in S^+$ that are of interest for Markman's results.

3.2. Representations of $spin(V)_s = so(7)$

For $s \in S^+$, the Cayley class c_s of s is the image of $s \odot s$ under the composition $Sym^2(S^+) \to \Gamma_{2\alpha} \hookrightarrow \wedge^4 V$. We will now assume that $(s, s)_{S+} \neq 0$. Then the stabilizer of s in $SO(S^+)$ is the orthogonal group $SO(s^\perp) \cong SO(7)$. The inverse image of this group in $Spin(V)$, a double cover of $SO(S^+)$, is denoted by $Spin(V)_s \cong Spin(7)$. In the standard representation ρ_V of $Spin(V)$ on V, the subgroup $Spin(V)_s$ still acts irreducibly, in fact V is isomorphic with the (unique, irreducible) spin representation of $Spin(7)$.

Since s is fixed by $Spin(V)_s$, the 4-form c_s is fixed by the Lie algebra $spin(V)_s \cong so(7)$. We now show that c_s is the unique $spin(V)_s$-invariant in $\wedge^4 V$ by considering the restriction to $so(7)$ of the $so(V) = so(8)$-representations appearing in Section 2.3.

Multiplication by s gives an inclusion of $spin(V)_s$-representations

$$
\begin{aligned}
S^+ = \langle s \rangle \oplus s^\perp \hookrightarrow Sym^2(S^+) &= \Gamma_0 \oplus \Gamma_{2\alpha} \\
&= \Gamma_0 \oplus \langle c_s \rangle \oplus s \odot s^\perp \oplus \Gamma_{(2,0,0)} \\
&= \Gamma_0 \oplus \Gamma_{(0,0,0)} \oplus \Gamma_{(1,0,0)} \oplus \Gamma_{(2,0,0)},
\end{aligned}
$$

where Γ_0 and $\Gamma_{(0,0,0)}$ are trivial $spin(V)_s$-representations, $\Gamma_{(1,0,0)} \cong s \odot s^\perp \cong s^\perp$ is the standard seven dimensional representation of $spin(V)_s \cong so(7)$ and $\Gamma_{(2,0,0)}$ is irreducible of dimension $35 - 1 - 7 = 27$ (the notation $\Gamma_{(a,b,c)}$ for $so(7)$-representations is as in [5]).

The restriction of the $spin(V)$-representation $\Gamma_{2\alpha}$ is thus a direct sum of three irreducible representations of $spin(V)_s$. The representation of $spin(V)_s$ on the other irreducible component $\Gamma_{2\beta}$ of $\wedge^4 V$ is irreducible and it is isomorphic to $\Gamma_{(0,0,2)}$. Thus one has the $spin(7) = so(7)$-decomposition into irreducible representations (cf. [14, Prop. 2], [7, Prop. 10.5.4]):

$$\wedge^4 V = \Gamma_{(0,0,0)} \oplus \Gamma_{(1,0,0)} \oplus \Gamma_{(2,0,0)} \oplus \Gamma_{(0,0,2)}.$$

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Since there is a unique copy of the trivial representation of $so(7)$ in $\wedge^4 V$, the Cayley class is the unique $spin(V)_s$ invariant in $\wedge^4 V$.

3.3.

The following proposition computes the 4-form c_s, which spans the trivial $spin(V)_s$-subrepresentation $\Gamma_{(0,0,0)}$ in $\wedge^4 V$, explicitly in a case of interest in Markman's paper, cf. [12, 1.4.1, Proposition 11.2]. There s is called $w = s_n$. We consider in fact $\frac{1}{n+1} c_w$ and we write n for his $n+1$. Notice that the computation below uses only representation theory.

3.4. Proposition

Let $n \in \mathbb{Z}$, $n \neq 0$, and let $s = s_n = 1 - ne_* \in S^+$. where $e_* := e_1 \wedge e_2 \wedge e_3 \wedge e_4 \in \wedge^{even} W = S^+$. Then we have, up to a scalar multiple,

$$c_s = -n\alpha^2 + 4n^2\beta + 4\gamma \quad (\in \wedge^4 V),$$

where the forms, now in $\wedge^* V$, involved are:

$$\alpha := e_1 \wedge e_5 + \cdots + e_4 \wedge e_8, \qquad \beta := e_1 \wedge \cdots \wedge e_4, \qquad \gamma := e_5 \wedge \cdots \wedge e_8.$$

Proof. The space of $spin(V)_s$-invariants in $\wedge^4 V$ is one dimensional and it is spanned by c_s, see Section 3.2. So it suffices to show that the right hand side is a non-zero $spin(V)_s$-invariant form.

The Lie algebra $spin(V)_{1,e_*}$ that acts trivially on the two dimensional subspace of S^+ spanned by $1, e_*$ is isomorphic to $so(6) \cong sl(4)$. The representation of $sl(4)$ on $V = W \oplus W^*$ is reducible and W is the standard representation of $sl(4)$ whereas W^* is the dual of the standard representation. This implies that $\beta \in \wedge^4 W \subset \wedge^4 V$ and $\gamma \in \wedge^4 W^* \subset \wedge^4 V$ as well as the 2-form α, which is the $sl(4)$-invariant in $W \otimes W^* \subset \wedge^2 V$ corresponding to the symplectic form $((w_1, w_1^*), (w_2, w_2^*)) = w_1^*(w_2) - w_2^*(w_1)$ on V, are $spin(V)_{1,e_*}$-invariants. On the other hand,

$$\wedge^4(W \oplus W^*) = \wedge^4 W \oplus W \otimes \wedge^3 W^* \oplus \wedge^2 W \otimes \wedge^2 W^* \oplus \wedge^3 W \otimes W^* \oplus \wedge^4 W^*.$$

Since W, W^* have dimension four, $\wedge^3 W^* \cong W$ and it is well-known that there are no $sl(4)$-invariants in $W \otimes W$ nor in $W^* \otimes W^*$. Also $\wedge^2 W$ is irreducible and thus the $sl(4)$-invariants in $\wedge^2 W \otimes \wedge^2 W^* \cong End(\wedge^2 W)$ are a one dimensional subspace spanned by the trace. Hence the subspace of $sl(4)$-invariants in $\wedge^4 V$ has dimension three. Since α^2, β, γ are linearly independent invariants, the invariant subspace is

$$(\wedge^4 V)^{spin(V)_{1,e_*}} = (\wedge^4 V)^{sl(4)} = \langle \alpha^2, \beta, \gamma \rangle.$$

Since $spin(V)_{1,e_*} \subset spin(V)_s$, any $spin(V)_s$-invariant in $\wedge^4 V$ must lie in the subspace $\langle \alpha^2, \beta, \gamma \rangle$. The 21-dimensional Lie algebra $spin(V)_s$ is defined by

$$spin(V)_s = \{ X \in spin(V) : Xs = 0 \}.$$

The action of $spin(V)$ on S^+ is given in the proof of Theorem 1.9. It is then easy to check that the following elements (of $so(2n) \cong spin(V)$) span $spin(V)_s$:

$$\mathfrak{h}_s := \left\{ \sum a_i X_{i,i} : \sum a_i = 0 \right\}, \quad X_{i,j} \quad (i \neq j), \quad nY_{i,j} \pm Z_{k,l} \quad (\{i, j, k, l\} = \{1, \ldots, 4\}),$$

where the sign depends on $i, \ldots, l$. In particular, $X := nY_{1,2} + Z_{3,4} \in spin(V)_s$ (in fact X acts as $e_1 e_2 + D_{e_3} D_{e_4}$ on S^+ and $X(1) = ne_1 e_2$, $X(e_*) = -e_1 e_2$, so $Xs = 0$). The action of X on V is given by

$$
\begin{array}{llllllll}
X(e_1) & = & 0, & X(e_2) & = & 0, & X(e_3) & = & -e_8, & X(e_4) & = & e_7, \\
X(e_5) & = & -ne_7, & X(e_6) & = & ne_8, & X(e_7) & = & 0, & X(e_8) & = & 0.
\end{array}
$$

Since the Lie algebra element X acts a derivation on $\wedge^4 V$ we have

$$
X(\alpha) = X(e_1) \wedge e_5 + e_1 \wedge X(e_5) + \cdots = -2ne_1 \wedge e_2 + 2e_7 \wedge e_8.
$$

Thus

$$
X(\alpha^2) = 2\alpha \wedge X(\alpha) = -4n(e_1 \wedge e_2) \wedge (e_3 \wedge e_7 + e_4 \wedge e_8) + 4(e_1 \wedge e_5 + e_2 \wedge e_6) \wedge (e_7 \wedge e_8).
$$

Similarly one finds

$$
X(\beta) = (e_1 \wedge e_2) \wedge (e_4 \wedge e_8 + e_3 \wedge e_7), \qquad X(\gamma) = -n(e_2 \wedge e_6 + e_1 \wedge e_5) \wedge (e_7 \wedge e_8).
$$

Therefore the only non-trivial linear combination of α^2, β, γ that is mapped to zero by X is $-n\alpha^2 + 4n^2\beta + 4\gamma$. Hence this must be the unique $spin(V)_s$-invariant in $\wedge^4 V$. $\qquad \square$

4. Abelian varieties of Weil type

4.1. The complex tori $\mathcal{T}_\ell$ and abelian varieties

For a point $\ell \in \Omega$, an open subset of the spinor quadric Q^+, we defined a complex torus $\mathcal{T}_\ell$ of dimension four whose first cohomology group is identified with V and whose Hodge structure is determined by $H^{1,0}(\mathcal{T}_\ell) = Z_\ell$, the maximal isotropic subspace of $V_{\mathbb{C}}$ corresponding to ℓ.

Fixing an $s \in S^+$ we also found that for $\ell \in \Omega_{s^\perp}$ this complex torus has an integral Hodge class (the Cayley class) $c_s \in H^{2,2}(\mathcal{T}_\ell, \mathbb{Z})$. Now we assume that $(s, s)_{S^+} > 0$ and we fix another, non-isotropic, class $h \in s^\perp$ with $(h, h)_{S^+} > 0$. Hence the rank two sublattice $\langle h, s \rangle \subset S^+$ generated by h, s is positive definite for the bilinear form on S^+. For $\ell \in \langle h, s \rangle^\perp \cap \Omega$, the torus T_ℓ turns out to be an abelian variety of Weil type and the Cayley class c_s is a non-trivial Hodge class. This result, Theorem 4.6, is due to O'Grady [15, Theorem 5.1] and Markman [12, Corollary 12.9, Theorem 13.4]. First we recall the basic facts on abelian varieties of Weil type.

4.2. Abelian varieties of Weil type

Let A be an abelian variety and let $K = \mathbb{Q}(\sqrt{-d})$, with $d \in \mathbb{Z}_{>0}$, be an imaginary quadratic field. An abelian variety of Weil type (with field K) is a pair (A, K), where A is an abelian variety and $K \hookrightarrow End(A)_{\mathbb{Q}}$ is a subalgebra of the endomorphism algebra of A, such that for all $x \in K$, $x \notin \mathbb{Q}$, the endomorphism of $T_0 A$ defined by the differential of $x = a + b\sqrt{-d} \in K$, with $a, b \in \mathbb{Q}$, has eigenvalues $x = a + b\sqrt{-d}$ and $\bar{x} = a - b\sqrt{-d}$ with the same multiplicity. Equivalently, the eigenvalues of x^* on $H^{1,0}$ have the same multiplicity. In particular, if (A, K) is of Weil type, then $\dim A$ is even.

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Given an abelian variety of Weil type (A, K), there exists a polarization $\omega_K \in H^{1,1}(A, \mathbb{Z})$ on A such that for all $x \in K$ its pull-back is

$$x^* \omega_K = \mathrm{Nm}(x) \omega_K, \qquad \mathrm{Nm}(x) = x\bar{x},$$

where $\mathrm{Nm}(x)$ is the norm of $x \in K$ (see [17, Lemma 5.2.1]). We call such a 2-form a polarization of Weil type and (A, K, ω_K) is called a polarized abelian variety of Weil type.

4.3. The Weil classes

For a general abelian variety of Weil type (A, K) of dimension $2n$, the spaces of Hodge classes

$$B^p(A) := H^{p,p}(A, \mathbb{Q}) := H^{2p}(A, \mathbb{Q}) \cap H^{p,p}(A)$$

have dimensions [21], see also [17, Theorem 6.12]:

$$\dim B^p(A) = 1, \quad (p \neq n), \qquad \dim B^n(A) = 3.$$

Since $\dim B^1(A) = 1$, any $\omega \in B^1(A)$, $\omega \neq 0$, defines (up to sign) a polarization on A which will be of Weil type.

The action of the multiplicative group $K^\times := K - \{0\}$ on $H^1(A, K) := H^1(A, \mathbb{Q}) \otimes_{\mathbb{Q}} K$ has an eigenspace decomposition into two $2n$-dimensional K subspaces

$$H^1(A, K) = Z_\kappa \oplus Z_{\bar\kappa}, \qquad x^*(v, w) = (xv, \bar{x}w)$$

that are conjugate over K. Since A is of Weil type, the complexifications of these eigenspaces both have Hodge numbers $h^{1,0} = h^{0,1} = n$. Thus in $H^{2n}(A, K) = \wedge^{2n} H^1(A, K)$ there are two 1-dimensional K-subspaces $\wedge^{2n} Z_\kappa$, $\wedge^{2n} Z_{\bar\kappa}$ of Hodge type (n, n). Since they are conjugate, their direct sum is defined over $\mathbb{Q}$, that is, there is a 2-dimensional $\mathbb{Q}$-subspace W_K of Hodge classes

$$W_K \subset H^{n,n}(A, \mathbb{Q}), \qquad W_K \otimes_{\mathbb{Q}} K = \wedge^{2n} Z_\kappa \oplus \wedge^{2n} Z_{\bar\kappa}.$$

(There is also a natural identification of W_K with $\wedge_K^{2n} H^1(A, \mathbb{Q})$ where $H^1(A, \mathbb{Q})$ is viewed as a $2n$-dimensional K vector space.) The subspace W_K is called the space of Weil classes. For any A of Weil type one has

$$\mathbb{Q}\omega_K^n \oplus W_K \subseteq B^n(A)$$

where ω_K^n is the n-fold exterior product of ω_K with itself. For a general A of Weil type one has $B^n(A) = \mathbb{Q}\omega_K^n \oplus W_K$.

An element $x \in K$ acts with eigenvalues $(x\bar{x})^n$, x^{2n}, $\bar{x}^{2n}$ on $\mathbb{Q}\omega_K^n \oplus W_K$. Thus if a non-zero element c in the three dimensional $\mathbb{Q}$ vector space $\mathbb{Q}\omega_K^n \oplus W_K$ is algebraic and it is not an eigenvector for the K-action (so it is not a multiple of ω_K^n) then all classes in $\mathbb{Q}\omega_K^n \oplus W_K$ are algebraic since ω_K^n is and so is $x^* c$ for all $x \in K$.

4.4. The Hermitian form

The $\mathbb{Q}$ vector space $H_1(A, \mathbb{Q})$ is also a K vector space for the action of K given by x_* for $x \in K \subset End(A)_{\mathbb{Q}}$. A polarization of Weil type $\omega_K \in H^2(A, \mathbb{Q})$ defines an

alternating form on $H_1(A, \mathbb{Q})$ and it also defines a K-valued Hermitian form H on this K-vector space by:

$$H : H_1(A, \mathbb{Q}) \times H_1(A, \mathbb{Q}) \longrightarrow K,$$
$$H(x, y) := \omega_K(x, (\sqrt{-d})_* y) + \sqrt{-d}\,\omega_K(x, y).$$

If $\Psi \in M_n(K)$ is the Hermitian matrix defining H w.r.t. some K-basis of $H_1(A, \mathbb{Q})$ then $\det(\Psi) \in \mathbb{Q}^\times = \mathbb{Q} - \{0\}$ and the class of $\det(\Psi) \in \mathbb{Q}^\times/\mathrm{Nm}(K^\times)$, called the discriminant of H, is independent of the choice of the basis. Given two non-degenerate Hermitian forms H, H' on K^n, there is a K-linear map $M : K^n \to K^n$ such that $H'(x, y) = H(Mx, My)$ for all $x, y \in K^n$ if and only if H, H' have the same signature and the same discriminant.

The discriminant of a polarized abelian variety of Weil type (A, K, ω_K) is the discriminant of H.

In Markman's approach, the real part of H, which is a bilinear form, is (up to the duality between $H_1(A, \mathbb{Z})$ and $H^1(A, \mathbb{Z})$ and up to a scalar multiple) the bilinear form $(\cdot, \cdot)_V$, cf. Section 4.8. In particular, it is the same for all families of Weil type, for all fields, considered in [12] and in Theorem 4.6 below.

4.5. Complete families

Given a K vector space U of dimension $2n$ and a Hermitian form $H : U \times U \to K$, any $2n$-dimensional abelian variety of Weil type A with field K and discriminant equal to the discriminant of H is obtained by choosing a free $\mathbb{Z}$-module $\Lambda \subset U$ of rank $4n$ and a complex structure J on $\Lambda_\mathbb{R} := \Lambda \otimes_\mathbb{Z} \mathbb{R}$ such that J commutes with K, the two eigenspaces of $x \in K$, $x \notin \mathbb{Q}$, on $(\Lambda_\mathbb{R}, J)$ have the same dimension and finally the imaginary part ω_K of H defines a polarization on the complex torus $(\Lambda_\mathbb{R}, J)/\Lambda$.

The unitary group $U(H)_\mathbb{R} \cong U(n, n)$ of the Hermitian form H on the $\mathbb{C} = K \otimes_\mathbb{Q} \mathbb{R}$ vector space $\Lambda_\mathbb{R}$ acts by conjugation $g \cdot J := gJg^{-1}$ on these complex structures. From this one obtains a complete family of abelian $2n$-folds of Weil type parametrized by a Hermitian symmetric domain isomorphic to $U(n, n)/(U(n) \times U(n))$, so of complex dimension n^2. The unitary group $SU(H) \subset GL(\Lambda_\mathbb{Q})$, viewed as algebraic group over $\mathbb{Q}$, is the special Mumford Tate group of the general abelian variety in the family, see [17].

We discuss the proof of the following theorem in the remainder of this section.

4.6. Theorem

Let $h, s \in S^+$ be perpendicular and such that $\langle h, s \rangle \subset S^+$ is a positive definite rank two sublattice. Let $d := (h, h)_{S^+}(s, s)_{S^+} \in \mathbb{Q}_{>0}$ and let $\ell \in \Omega_{\{h,s\}^\perp}$, where

$$\Omega_{\{h,s\}^\perp} := \{\ell \in \Omega_{s^\perp} : \quad (\ell, h)_{S^+} = 0\} = \{\ell \in \Omega : \quad (\ell, s)_{S^+} = (\ell, h)_{S^+} = 0\}$$

is a complex manifold of dimension four. Then we have:

(a) The complex four dimensional torus $\mathcal{T}_\ell$ has endomorphisms by $K = \mathbb{Q}(\sqrt{-d})$, that is $K \subset \mathrm{End}(\mathcal{T}_\ell)_\mathbb{Q}$.
(b) The complex torus $\mathcal{T}_\ell$ has a polarization $\omega_K \in H^2(\mathcal{T}_\ell, \mathbb{Z})$ and $(\mathcal{T}_\ell, K, \omega_K)$ is polarized abelian fourfold of Weil type.

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(c) The discriminant of the polarization $\omega_K \in H^2(\mathcal{T}_\ell, \mathbb{Z})$ is trivial.
(d) The Cayley class $c_s \in H^{2,2}(\mathcal{T}_\ell, \mathbb{Z})$ is not contained in the subspace $\mathbb{Q}\omega_K^2$ where $\omega_K^2 = \omega_K \wedge \omega_K$.
(e) The four dimensional family of these fourfolds of Weil type parametrized by $\Omega_{\{h,s\}^\perp}$ is complete.

4.7. Endomorphisms of $\mathcal{T}_\ell$

Since the sublattice $\langle h, s \rangle$ is positive definite, we may assume that the restriction q of the quadratic form on S^+ is given by $q(xh + ys) = ax^2 + by^2$, with both $a = (h, h)_{S^+}, b = (s, s)_{S^+} \in \mathbb{Q}$ positive. Hence $d = ab > 0$. The zero locus of q is defined by $a^{-1}((ax)^2 + aby^2) = 0$, showing that there are two isotropic lines in $\langle h, s \rangle_{\mathbb{C}}$ defined by $ax \pm \sqrt{-d}y = 0$. These two lines are conjugate over K where the conjugation on K is $x + y\sqrt{-d} = x - y\sqrt{-d}$ with $x, y \in \mathbb{Q}$. In $\mathbf{P}S_{\mathbb{C}}^+$ they correspond to the two points of intersection of the line $\mathbf{P}\langle h, s \rangle_{\mathbb{C}}$ with the spinor quadric $Q^+ \cong IG(4, V_{\mathbb{C}})^+$, which we denote by $\kappa, \overline{\kappa}$:

$$\{\kappa, \overline{\kappa}\} = Q^+ \cap \mathbf{P}\langle h, s \rangle_{\mathbb{C}} \qquad (\subset \mathbf{P}S_{\mathbb{C}}^+).$$

As $Q^+ = \gamma(IG(4, V_{\mathbb{C}})^+)$, these two points define two maximal isotropic subspaces in $V_K := V \otimes_{\mathbb{Q}} K$ denoted by $Z_\kappa, Z_{\overline{\kappa}}$. Since the points $\kappa, \overline{\kappa}$ are conjugate over K, so are these subspaces: if $w = v + \sqrt{-d}v' \in Z_\kappa$ with $v, v' \in V_{\mathbb{Q}}$ then $\overline{w} = v - \sqrt{-d}v' \in Z_{\overline{\kappa}}$.

The plane $\langle h, s \rangle_{\mathbb{C}}$ is not contained in Q^+, hence these two subspaces have trivial intersection (Lemma 1.11, [3, III.1.12]):

$$V_K = Z_\kappa \oplus Z_{\overline{\kappa}}, \qquad \overline{(v_1, v_2)} = (\overline{v_2}, \overline{v_1}) \qquad (v_1 \in Z_\kappa, \ v_2 \in Z_{\overline{\kappa}}).$$

We identify the $\mathbb{Q}$ vector space $V_{\mathbb{Q}}$ with the image of $V_{\mathbb{Q}} \hookrightarrow V_K$, it consists of the points $(v_1, \overline{v_1})$ with $v_1 \in Z_\kappa$. Now we define an action of K on $V_{\mathbb{Q}}(\subset V_K)$ by

$$K \times V_{\mathbb{Q}} \longrightarrow V_{\mathbb{Q}}, \qquad x \cdot (v_1, \overline{v_1}) := (xv_1, \bar{x}\overline{v_1}) = (xv_1, \overline{xv_1}) \qquad (\in V_{\mathbb{Q}} \subset Z_\kappa \oplus Z_{\overline{\kappa}}),$$

where $\bar{x}$ is the conjugate of $x \in K$.

To show that this induces an inclusion $K \subset \mathrm{End}(\mathcal{T}_\ell)_{\mathbb{Q}}$, it suffices to verify that any $x \in K$ commutes with the complex structure J_ℓ on $V_{\mathbb{R}}$. Since $\ell \in \Omega_{h,s^\perp}$ we have $(\ell, \kappa)_{S^+} = 0$ and similarly the scalar products of any one of $\ell, \overline{\ell}$ and any one of $\kappa, \overline{\kappa}$ are zero. Therefore the intersection of $Z_\ell, Z_{\overline{\ell}}$ with the complexifications of $Z_\kappa, Z_{\overline{\kappa}}$ is not zero by Lemma 1.11. Since these spaces are parametrized by the same connected component $IG(4, V_{\mathbb{C}})^+$, their intersection is even dimensional and thus it is two dimensional. From the eigenspace decomposition for J_ℓ, $V_{\mathbb{C}} = Z_\ell \oplus Z_{\overline{\ell}}$, we obtain the decomposition

$$V_{\mathbb{C}} = (Z_\ell \cap Z_{\kappa, \mathbb{C}}) \oplus (Z_\ell \cap Z_{\overline{\kappa}, \mathbb{C}}) \oplus (Z_{\overline{\ell}} \cap Z_{\kappa, \mathbb{C}}) \oplus (Z_{\overline{\ell}} \cap Z_{\overline{\kappa}, \mathbb{C}}).$$

The action of J_ℓ and $x \in K$ on these four summands are scalar multiplications (by $\pm i$ and $x, \bar{x}$ respectively), hence the action of K indeed commutes with J_ℓ. Since each summand has dimension 2, the eigenvalues of $x \in K$, $x \notin \mathbb{Q}$, on $Z_\ell = H^{1,0}(\mathcal{T}_\ell)$ have the same dimension.

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4.8. The polarization

The combination of the K-action on $V_{\mathbb{Q}} = H^1(\mathcal{T}_\ell, \mathbb{Q})$ with the bilinear form $(\bullet, \bullet)_V$ leads a polarization $\omega_K \in H^2(\mathcal{T}_\ell, \mathbb{Q})$ on $\mathcal{T}_\ell$. We define a bilinear form E on $V_{\mathbb{Q}}$ by:

$$E : V \times V \longrightarrow \mathbb{Q}, \qquad E(v, w) = (\sqrt{-d} \cdot v, w)_V.$$

The duality $V = H_1(\mathcal{T}_\ell, \mathbb{Q})^{dual}$ implies that E defines an element $\omega_K \in \wedge^2 V = H^2(\mathcal{T}_\ell, \mathbb{Q})$. Similar to the computations for Kähler forms and metrics we establish the basic properties of E which imply that $(\mathcal{T}_\ell, K, \omega_K)$ is a polarized abelian fourfold of Weil type.

First of all, we have for all $v, w \in V_{\mathbb{Q}}$ and all $x \in K$ that

$$E(x \cdot v, x \cdot w) = x\bar{x} E(v, w).$$

To verify this, we extend E K-bilinearly to V_K and we use that $Z_\kappa, Z_{\bar\kappa}$ are isotropic subspaces. Thus, with $v = v_1 + \overline{v_1}$, $w = w_1 + \overline{w_1} \in Z_\kappa \oplus Z_{\bar\kappa}$ we get

$$
\begin{aligned}
E(x \cdot v, x \cdot w) &= \left(x\sqrt{-d}v_1 + \overline{x\sqrt{-d}v_1}, xw_1 + \overline{xw_1}\right)_V \\
&= \left(x\sqrt{-d}v_1, \bar{x}\overline{w_1}\right)_V + \left(\bar{x}\sqrt{-d}v_1, xw_1\right)_V \\
&= x\bar{x}\left((\sqrt{-d}v_1, \overline{w_1})_V + (\sqrt{-d}v_1, w_1)_V\right) \\
&= x\bar{x} E(v, w).
\end{aligned}
$$

Next we show that E is alternating:

$$E(v, w) = (\sqrt{-d} \cdot v, w)_V = (w, \sqrt{-d} \cdot v)_V = \frac{1}{d}(\sqrt{-d} \cdot v, \sqrt{-d}^2 \cdot w)_V$$
$$= -(\sqrt{-d} \cdot v, w)_V = -E(w, v).$$

To show that the 2-form ω_K is of type $(1, 1)$ it suffices to show that $E(J_\ell v, J_\ell w) = E(v, w)$ for all $v, w \in V_{\mathbb{R}}$:

$$E(J_\ell v, J_\ell w) = \left(\sqrt{-d} \cdot J_\ell v, J_\ell w\right)_V = \left(J_\ell(\sqrt{-d} \cdot v), J_\ell w\right)_V = \left(\sqrt{-d} \cdot v, w\right)_V = E(v, w).$$

Finally we verify that $E(J_\ell v, v) > 0$ for non-zero $v \in V_{\mathbb{R}}$. That is, we must show that $(\sqrt{-d} \cdot J_\ell v, w) > 0$. The endomorphisms $\sqrt{-d}, J_\ell$ of $V_{\mathbb{R}}$ are both constructed from decompositions of $V_{\mathbb{C}}$ with two conjugate isotropic subspaces $Z_\kappa, Z_{\bar\kappa}$ and $Z_\ell, Z_{\bar\ell}$ respectively. The corresponding points $\kappa, \bar\kappa, \ell, \bar\ell \in Q^+ = IG(4, V)^+$ span a $\mathbf{P}^3 \in \mathbf{P}S_{\mathbb{C}}^+$ which is the projectivization of the complexification of the four dimensional subspace $< h, s, \ell + \bar\ell, (\ell - \bar\ell)/i > \subset S_{\mathbb{R}}^+$ (here $\mathbb{C} = \mathbb{R} + i\mathbb{R}$). Notice that this basis consists of perpendicular vectors for $(\bullet, \bullet)_{S^+}$ and that the subspace is positive definite.

The group $Spin(V_{\mathbb{R}})$ acts via $SO(S_{\mathbb{R}}^+)$ on $S_{\mathbb{R}}^+$ and this action is transitive on such subspaces. As $Spin(V_{\mathbb{R}})$ also acts via $SO(V_{\mathbb{R}})$ on $V_{\mathbb{R}}$, we see that it suffices to show that $(J_1 J_2 v, v) > 0$ for all non-zero $v \in V_{\mathbb{R}}$ where the linear maps J_1, J_2 are defined by any two orthogonal positive definite 2-dimensional subspaces of $S_{\mathbb{R}}^+$. (Markman shows that the map $J_1 J_2$ is already determined, up to a scalar multiple, by the direct sum of these subspaces.)

We use the conventions from Section 1.10. A point $z = (z_1, \ldots, z_8) \in S^+ \cong U^4$ will be written as

$$z = \left(\begin{pmatrix} z_1 \\ z_5 \end{pmatrix}, \begin{pmatrix} z_2 \\ z_6 \end{pmatrix}, \begin{pmatrix} z_3 \\ z_7 \end{pmatrix}, \begin{pmatrix} z_4 \\ z_8 \end{pmatrix}\right), \qquad (z, z)_{S^+} = 2(z_1 z_5 + \cdots + z_4 z_8).$$

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The following four points $v_1, \ldots, v_4$, where $v := \binom{1}{1} \in U$, in S^+ are perpendicular and span a positive 4-plane in $S^+_{\mathbb{R}}$ since $(v_i, v_i)_{S^+} = 8$ and we also define $\ell_1, \ell_2 \in S^+_{\mathbb{C}}$:

$$
\begin{aligned}
v_1 &= (v, \ v, \ v, \ v), \\
v_2 &= (v, \ v, -v, -v), \\
v_3 &= (v, -v, \ v, -v), \\
v_4 &= (v, -v, -v, \ v),
\end{aligned}
\qquad
\begin{aligned}
\ell_1 &:= (v_1 + iv_2)/(1+i) = (v, \ v, -iv, -iv), \\[4pt]
\ell_2 &:= (v_3 + iv_4)/(1+i) = (v, -v, -iv, \ iv).
\end{aligned}
$$

Then $\ell_1, \overline{\ell_1}$ and $\ell_2, \overline{\ell_2}$ are all isotropic vectors and they span $\langle v_1, v_2 \rangle_{\mathbb{C}}$ and $\langle v_3, v_4 \rangle_{\mathbb{C}}$ respectively. Isotropic vectors are in $Q^+ = \gamma(IG(4, V_{\mathbb{C}})^+)$ and since these four all have first coordinate $z_1 = 1$ they are in the image of the open set $IG(4, V_{\mathbb{C}})^+_0$ parametrized by the alternating 4×4 matrices. Using the explicit description of γ one finds

$$
\ell_k = \gamma(Z_{B_k}) \quad (k = 1, 2), \qquad
B_1 = \begin{pmatrix} 0 & 1 & -i & -i \\ -1 & 0 & i & -i \\ i & -i & 0 & -1 \\ i & i & 1 & 0 \end{pmatrix},
$$

$$
B_2 = \begin{pmatrix} 0 & -1 & -i & i \\ 1 & 0 & -i & -i \\ i & i & 0 & 1 \\ -i & i & -1 & 0 \end{pmatrix}.
$$

The eigenspace with eigenvalue $-1 = i^2 = (-i)^2$ of the endomorphism $J_1 J_2$ of $V_{\mathbb{R}}$ is the direct sum of $Z_{\ell_1} \cap Z_{l_2}$ and its complex conjugate. Let c_k, d_k denote the kth column of the matrix $\binom{B_1}{I}$, $\binom{B_2}{I}$ respectively, then Z_{ℓ_1}, Z_{ℓ_2} are spanned by the c_k and the d_k $(k = 1, \ldots, 4)$ respectively. Their intersection is spanned by

$$
c_1 - ic_3 = d_1 - id_3, \quad c_2 - ic_4 = d_2 - id_4 \qquad (\in Z_{\ell_1} \cap Z_{\ell_2}).
$$

Considering $(c_1 - ic_3) \pm \overline{(c_1 - ic_3)}$ etc., one finds a basis of the -1-eigenspace of J_1, J_2. Its perpendicular is the $+1$-eigenspace. Recall that $e_1, \ldots, e_8$ are the basis vectors of V as in 0.1, then the eigenspace decomposition is:

$$
V_{\mathbb{R}} = V_+ \oplus V_- = \langle e_1 + e_5, e_2 + e_6, e_3 + e_7, e_4 + e_8 \rangle_{\mathbb{R}} \oplus \langle e_1 - e_5, e_2 - e_6, e_3 - e_7, e_4 - e_8 \rangle_{\mathbb{R}}.
$$

Notice that $(\bullet, \bullet)_V$ is positive definite on V_+ and negative definite on V_-. Writing $v = v_+ + v_-$ as sum of $J_1 J_2$ eigenvectors, one has $(J_1 J_2 v, v)_V = (v_+, v_+)_V - (v_-, v_-)_V$ and thus indeed $(J_1 J_2 v, v)_V > 0$ for all non-zero $v \in V_{\mathbb{R}}$.

4.9. The discriminant

We refer to [12, Lemma 12.11] (cf. [15, Theorem 5.1]) for the computation of the discriminant. See also Proposition 6.5 for a proof of the triviality of the discriminant using results from Lombardo [10].

4.10. The Cayley class and the Weil classes

We define two subgroups of $Spin(V)$. Let $Spin(V)_s$ be the subgroup which fixes $s \in S^+$ and let $Spin(V)_{h,s}$ be the subgroup which fixes all elements in $\langle h, s \rangle$. Then one

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can show that the Cayley class c_s is the unique $Spin(V)_s$-invariant in $\wedge^4 V$ and that ω_K is the unique $Spin(V)_{s,h}$-invariant in $\wedge^2 V$. This implies that $c_s \notin \mathbb{Q}\omega_K^2$ (cf. [14, Prop 2], [12, Thm 13.4] and Section 3.2).

One can also use that the $K \otimes_{\mathbb{Q}} \mathbb{C} \cong \mathbb{C} \times \mathbb{C}$-action on $V_{\mathbb{C}}$ has the eigenspaces $(Z_\kappa)_{\mathbb{C}}, (Z_{\bar{\kappa}})_{\mathbb{C}}$. The one parameter subgroup h_R of $Spin(V_{\mathbb{C}})$ which acts as multiplication by t, t^{-1} respectively on these eigenspaces fixes E, and thus it fixes $\omega_K \in \wedge^2 V$ and also $\omega_K^2 \in \wedge^4 V$. On the other hand, h_R has eigenvalues t^2, t^{-2} on $\langle \kappa, \bar{\kappa} \rangle_{\mathbb{C}} = \langle h, s \rangle_{\mathbb{C}} \subset S_{\mathbb{C}}^+$ by Lemma 2.6. Therefore c_s, the image of $s \odot s$ in $\wedge^4 V$, is not invariant under h_R and thus it cannot be a multiple of ω_K^2.

The Cayley class c_s is thus a Hodge class on any abelian variety of Weil type $\mathcal{T}_\ell$ with $\ell \in \Omega_{\{h,s\}^\perp}$, for any $h \in s^\perp$ such that $\langle h, s \rangle$ is positive definite. Such cohomology classes were found in a different context in [19], see Remark 5.3 of that paper for the relation with Markman's construction.

4.11. Complete families

The Lie group $Spin(V_{\mathbb{R}})_{h,s}$ acts on $\Omega_{\{h,s\}^\perp}$. This action induces an action of $Spin(V_{\mathbb{R}})_{h,s}$ on the orthogonal complex structures on $V_{\mathbb{R}}$ by $J_{g \cdot \ell} = g J_\ell g^{-1}$. The fixed points $\kappa, \bar{\kappa} \in Q^+ \cap \langle h, s \rangle_{\mathbb{C}}$ of the action of $Spin(V_{\mathbb{R}})_{h,s}$ on Q^+ correspond to the eigenspaces $Z_{\kappa,\mathbb{C}}, Z_{\bar{\kappa},\mathbb{C}}$ of the K-action, which are thus mapped into themselves. This implies that the image of $Spin(V_{\mathbb{R}})_{h,s}$ in $SO(V_{\mathbb{R}})$ commutes with the K action on $V_{\mathbb{R}}$. This image thus preserves the Hermitian form H and therefore $Spin(V)_{h,s}$ maps to the algebraic group $SU(H)$ which is the Mumford Tate group of the general $\mathcal{T}_\ell$ with $\ell \in \Omega_{\{h,s\}^\perp}$. For dimension reasons this map is surjective on the real points of these groups and thus the family of abelian fourfolds of Weil type is complete.

5. Moduli spaces of sheaves on an abelian surface

5.1.

The constructions considered thus far have a natural geometrical interpretation in terms of moduli spaces of sheaves on abelian surfaces. We now briefly recall the basic definitions and results, due to Mukai and Yoshioka. The notation used thus far is now adapted to this context, for example, the free $\mathbb{Z}$-module W of rank four will become $W = H^1(X, \mathbb{Z})$ for an abelian surface X etc.

We conclude with a brief outline of Markman's proof of the Hodge conjecture for the general abelian fourfolds of Weil type with trivial discriminant.

5.2. The Mukai lattice of an abelian surface

Let X be an abelian surface and let $\hat{X} = Pic^0(X)$ be the dual abelian surface. Let

$$W = H^1(X, \mathbb{Z}), \quad W^* = H^1(\hat{X}, \mathbb{Z}) = H^1(X, \mathbb{Z})^*, \qquad V := W \oplus W^*.$$

The Chern character of a coherent sheaf on X takes values in

$$S^+ := \wedge^{even} H^*(X, \mathbb{Z}) = H^0(X, \mathbb{Z}) \oplus H^2(X, \mathbb{Z}) \oplus H^4(X, \mathbb{Z}),$$

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and we will identify $H^0(X, \mathbb{Z})$, $H^4(X, \mathbb{Z})$ with $\mathbb{Z}$, using the generators 1 and a volume form compatible with the orientation on the complex manifold X.

The Mukai lattice of X is the (free, rank 8) $\mathbb{Z}$-module S^+ with the bilinear form given by (this bilinear form coincides up to sign with $(\bullet, \bullet)_{S^+}$):

$$(r, c, s) \cdot (r', c', s') := -(rs' + r's) + c \wedge c'.$$

For $v = (r, c, s) \in S^+$, with $r > 0$, $c \in NS(X) \subset H^2(X, \mathbb{Z})$ and $v^2 \geq 6$ the moduli space of sheaves E on X with $ch(E) = v$, denoted by $\mathcal{M}(v)$, is a smooth holomorphic symplectic manifold of dimension $v^2 + 2$.

5.3. The case $v = s_n$

We now take $v = s_n = (1, 0, -n)$, such that $v^2 = 2n \geq 6$ and $\dim M(v) = 2n + 2$. Let $Z \subset X$ be a subscheme of length n. Then its ideal sheaf $\mathcal{I}_Z$ has $ch(\mathcal{I}_Z) = v$ (for an abelian surface, the Chern character $ch(E)$ is the Mukai vector $v(E)$ of the sheaf E). This induces an inclusion of complex manifolds

$$Hilb^n(X) = X^{[n]} \hookrightarrow \mathcal{M}(v) \qquad (v = s_n = (1, 0, -n)).$$

For $\mathcal{L} \in \hat{X}$ and $\mathcal{I}_Z \in X^{[n]}$ one also has $\mathcal{L} \otimes \mathcal{I}_Z \in \mathcal{M}(v)$.

The Albanese map $\alpha : X^{[n]} \to X$ of $X^{[n]}$ fits in a diagram:

$$
\begin{array}{ccc}
X^{[n]} & & \\
\downarrow \quad \searrow \alpha & & \Sigma([p_1, \ldots, p_n]) := p_1 + \cdots + p_n, \\
X^{(n)} \xrightarrow{\Sigma} X & &
\end{array}
$$

here $X^{(n)}$ is the nth symmetric power of X and $[p_1, \ldots, p_n] \in X^{(n)}$ is the image of $(p_1, \ldots, p_n) \in X^n$ in $X^{(n)}$.

The generalized Kummer variety $K_{n-1}(X)$, of dimension $2n - 2$, is the irreducible holomorphic symplectic manifold obtained as

$$K_{n-1}(X) = \alpha^{-1}(0) \subset X^{[n]}.$$

Using locally free resolutions of sheaves one defines a determinant map $\det : M(v) \to \hat{X}$ and one has $\det(\mathcal{L} \otimes \mathcal{I}_Z) = \mathcal{L}$ for $\mathcal{L} \in \hat{X}$. Yoshioka [22] showed that

$$M(v) \cong \hat{X} \times (\det^{-1})(\mathcal{O}_X) \cong \hat{X} \times X^{[n]} \cong \hat{X} \times \left((X \times K_{n-1}(X))/X[n]\right)$$

where $X[n] \subset X$ is the subgroup of n-torsion points. In particular, the Bogomolov decomposition of $M(v)$ is the product of the abelian fourfold $X \times \hat{X}$ and the irreducible holomorphic symplectic manifold $K_{n-1}(X)$.

5.4. The cohomology of the generalized Kummer variety

The composition of the Mukai homomorphism [22, Section 1.2] and the restriction map

$$v^{\perp} \longrightarrow H^2(M(v), \mathbb{Z}) \longrightarrow H^2(K_{n-1}(X), \mathbb{Z})$$

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induces a Hodge isometry (for the weight two Hodge structure on $v^\perp$ defined by $(v^\perp)^{2,0} = H^{2,0}(X)$ and with the BBF quadratic form on $H^2(K_{n-1}(X), \mathbb{Z})$) [22, Thm. 0.2].

This implies, by the surjectivity of the period map and with $v = s_n = s$, that $\Omega_{s^\perp}$ is the period space of deformations of $K_{n-1}(X)$, these deformations are called Kummer type varieties.

Moreover, $h^{3,0}(K_{n-1}(X)) = 0$ so that $H^3(K_{n-1}(X), \mathbb{C}) = H^{2,1} \oplus H^{1,2}$ is essentially the first cohomology group of its intermediate Jacobian $H^3(\mathbb{C})/(H^{2,1} \oplus H^3(\mathbb{Z}))$ and one has ([22, Prop. 4.20]):

$$H^3(K_{n-1}(X), \mathbb{Z}) = H^1(X, \mathbb{Z}) \oplus H^3(X, \mathbb{Z}) \cong H^1(X, \mathbb{Z}) \oplus H^1(\hat{X}, \mathbb{Z}) = V.$$

O'Grady and Markman showed that for $\ell \in \Omega_{s^\perp}$ and any deformation Y_ℓ of $K_{n-1}(X)$ with period $H^{2,0}(Y_\ell) = \mathbb{C}\ell \subset (s^\perp)_\mathbb{C}$, there is an isomorphism of Hodge structures (up to Tate twist and isogeny) $H^3(Y_\ell, \mathbb{Z}) = H^1(\mathcal{T}_\ell, \mathbb{Z})$. In case the complex manifold Y_ℓ is algebraic and $h \in H^2(Y, \mathbb{Z}) = s^\perp$ is the class of an ample divisor, hence $\ell \in \Omega_{\{h,s\}^\perp}$, O'Grady [15] showed that the torus $\mathcal{T}_\ell$ is an abelian variety of Weil type. Moreover, he showed that for algebraic Y_ℓ the Kuga Satake variety of the weight two polarized Hodge structure of rank six $h^\perp \subset H^2(Y_\ell, \mathbb{Z})$ is (isogeneous to) $\mathcal{T}_\ell^4$ (see also Section 6.4 where $h^\perp \cong H_\ell^2$).

O'Grady also makes a detailed study of the cohomology of generalized Kummer varieties and in particular he observes that there is a natural map (recall $\dim Y_\ell = \dim K_{n-1}(X) = 2n - 2$):

$$H^3(Y_\ell, \mathbb{Z}) \longrightarrow H^{4n-6}(Y_\ell, \mathbb{Z}) \longrightarrow H^2(Y_\ell, \mathbb{Z})^\vee,$$

the last map is Poincaré duality, which relates the Hodge structures on $H^3(Y_\ell)$ and $H^2(Y_\ell)$.

5.5. Markman's theorem

Given a sheaf $F' \in \mathcal{M}(v)$ ($v = s_n$ as in Section 5.3), there is a natural map

$$\iota_{F'} : X \times \hat{X} \longrightarrow \mathcal{M}(v), \qquad (x, \mathcal{L}) \longmapsto (t_x^* F') \otimes \mathcal{L}$$

where $t_x : X \to X$, $y \mapsto x + y$ is the translation by x. Deforming $K_{n-1}(X)$ to Y_ℓ, with $\ell \in \Omega_{s^\perp}$, this map deforms to a map

$$\iota : \mathcal{T}_\ell \longrightarrow Y_\ell.$$

A universal sheaf $\mathcal{E}$ on $X \times \mathcal{M}(v)$ defines a sheaf E on $M(v) \times M(v)$ by $E := \mathcal{E}xt_{\pi_{13}}^1(\pi_{12}^* \mathcal{E}, \pi_{23}^* \mathcal{E})$ where π_{ij} are the projections from $M(v) \times X \times M(v)$. For $F \in M(v)$ let E_F the restriction of E to $\{F\} \times M(v) = M(v)$. This defines a sheaf on $X \times \hat{X}$ whose second Chern class is exactly the Cayley class defined by $v = s_n \in S^+$ ([12, Prop. 11.2], see also Proposition 3.4):

$$c_2(\iota_{F'}^* \mathcal{E}nd(E_F)) = c_v \in \wedge^4 V = H^4(X \times \hat{X}, \mathbb{Z}).$$

Markman, using results of Verbitsky, shows that the sheaf E_F on $M(v)$ deforms to a sheaf over any deformation Y_ℓ of $M(v)$. Thus $c_v \in H^4(\mathcal{T}_\ell, \mathbb{Z})$ is an algebraic class whenever $\mathcal{T}_\ell$ is an abelian variety. From Theorem 4.6.d we have that c_v is not

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an eigenvector for the action of the multiplicative group $K^\times$ on the Hodge classes in $\mathbb{Q}\omega_K^2 \oplus W_K \subset H^{2,2}(\mathcal{T}_\ell, \mathbb{Z})$. Thus ω_K^2, c_v and the images of c_v under the $K^\times$ action span $\mathbb{Q}\omega_K^2 \oplus W_K$. Since any fourfold of Weil type with trivial discriminant is isogeneous to a $\mathcal{T}_\ell$, for any such fourfold the space W_K is spanned by algebraic classes.

6. Kuga Satake varieties and abelian fourfolds of Weil type

6.1. Kuga Satake varieties

Let S^+ be the lattice introduced in Section 1.10. As in Theorem 4.6, let $h, s \in S^+ \cong U^{\oplus 4}$ be two perpendicular elements such that their span is a positive definite sublattice. Then for any $\ell \in \Omega_{\{h,s\}^\perp}$ the complex torus $\mathcal{T}_\ell$ is an abelian variety of Weil type. We now define a polarized weight two Hodge structure $H = H_\ell$ of rank six with Hodge numbers $\dim H_\ell^{2,0} = 1$, $\dim H_\ell^{1,1} = 4$. The Kuga Satake construction associates to H_ℓ an abelian variety A_ℓ of dimension $2^{6-1} = 32$. We show here, using results of Lombardo [10] and O'Grady [15], that A_ℓ is isogeneous to $\mathcal{T}_\ell^4$. Moreover, we provide an alternative proof of the fact that discriminant of the polarization of $\mathcal{T}_\ell$ is trivial.

Let $H = H_{h,s}$ be the rank 6 sublattice of signature $(2+, 4-)$ orthogonal to $\langle h, s \rangle$:

$$H := \langle h, s \rangle^\perp = \{ t \in S^+ : (t, h) = (t, s) = 0 \}.$$

With this notation we have

$$\Omega_{\{h,s\}^\perp} = \{ \ell \in \mathbf{P}H_\mathbb{C} : \quad (\ell, \ell)_{S^+} = 0, \quad (\ell, \bar\ell)_{S^+} > 0 \}.$$

Recall that any $\ell \in \Omega_{\{h,s\}^\perp}$ defines an abelian fourfold of Weil type with underlying torus $\mathcal{T}_\ell$ by Theorem 4.6. Such an ℓ also defines a weight two Hodge structure on H denoted by H_ℓ as follows:

$$H_{\ell,\mathbb{C}} = H_\mathbb{C} = \oplus_{p+q=2} H_\ell^{p,q}, \qquad H_\ell^{2,0} := \mathbb{C}\ell, \qquad H_\ell^{0,2} := \mathbb{C}\bar\ell,$$
$$H_\ell^{1,1} = \left(H_\ell^{2,0} \oplus H_\ell^{0,2} \right)^\perp.$$

This Hodge structure is polarized since the restriction of $(\bullet, \bullet)_{S^+}$ to the two dimensional real subspace $(H_\ell^{2,0} \oplus H_\ell^{0,2}) \cap H_\mathbb{R}$ is positive definite.

As $\dim H_\ell^{2,0} = 1$, there is a Kuga Satake (abelian) variety A_ℓ, of dimension 16, associated to H_ℓ (see [4,8,18]). In general, it has the property that H_ℓ is a Hodge substructure of $H^2(A_\ell^2, \mathbb{Q})$, but in this case there are actually several copies of H_ℓ in $H^2(A_\ell, \mathbb{Q})$, see Section 6.4. The even Clifford algebra $C(H)^+$ of H is a lattice in the real vector space $C(H)^+ \otimes_\mathbb{Z} \mathbb{R}$ of dimension $2^5 = 32$. A complex structure on $C(H)_\mathbb{R}^+$ is defined by left multiplication by $f_1 f_2 \in C(H)_\mathbb{R}^+$, with $f_1, f_2 \in H_\mathbb{R}$ such that $(f_1, f_1)_{S^+} = 1$ and $H_\ell^{2,0} = \langle f_1 + if_2 \rangle$ (cf. [18, Section 5.6]). The abelian variety A_ℓ is the quotient $(C(H)_\mathbb{R}^+, f_1 f_2)/C(H)^+$.

In [10, Cor. 6.3, Thm 6.4] it is shown that A_ℓ is isogeneous to B_ℓ^4, where B_ℓ is an abelian fourfold of Weil type with trivial discriminant. The following proposition, due to O'Grady ([15, Section 5.3]), shows that B_ℓ and $\mathcal{T}_\ell$ are isogeneous. In [15] one finds a more explicit description of this result, as well as applications to generalized Kummer varieties.

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6.2. Proposition

For $\ell \in \Omega_{\{h,s\}^{\perp}}$ the Kuga Satake variety A_ℓ of the polarized weight two Hodge structure H_ℓ is isogeneous to $\mathcal{T}_\ell^4$, where $\mathcal{T}_\ell$ is the abelian fourfold of Weil type defined by ℓ.

Proof. The right multiplication on $C(H)_{\mathbb{R}}^+$ by an element of $C(H)^+$ preserves the lattice, commutes with the complex structure and thus defines an element in $\mathrm{End}(A_\ell)$. The $\mathbb{Q}$ vector space $H_{\mathbb{Q}}$ is not a direct sum of two maximally isotropic subspaces and, whereas $C(H)_{\mathbb{C}}^+ \cong M_4(\mathbb{C}) \oplus M_4(\mathbb{C})$ (as in Section 1.2), one now has an isomorphism of algebras ([10, Thm. 6.2]), where $M_4(K)$ are the 4×4 matrices with coefficients in K,

$$C(H)_{\mathbb{Q}}^+ := C(H)^+ \otimes_{\mathbb{Z}} \mathbb{Q} \cong M_4(K) \subseteq \mathrm{End}(A_\ell)_{\mathbb{Q}}, \qquad K := \mathbb{Q}(\sqrt{-ab}).$$

This implies that any A_ℓ is isogeneous to B_ℓ^4, where B_ℓ is an abelian fourfold with $K \subset \mathrm{End}(B_\ell)_{\mathbb{Q}}$ (B_ℓ is only determined up to isogeny).

It remains to show that B_ℓ and $\mathcal{T}_\ell$ are isogeneous. The inclusion $Spin(H) \subset Spin(S^+) = Spin(V)$ defines a representation of $Spin(H)$ on V which is its spin representation. The isomorphism $C(H)_{\mathbb{Q}}^+ \cong M_4(K)$ implies that

$$C(H)_{\mathbb{Q}}^+ \cong V_{\mathbb{Q}}^{\oplus 4}$$

as $Spin(H)$-representations. The same holds with $\mathbb{Q}$ replaced by $\mathbb{R}$. The weight two Hodge structure on the $Spin(H)$-representation H_ℓ is defined by the one parameter subgroup h_ℓ of $Spin(H)_{\mathbb{R}} \subset Spin(S^+)_{\mathbb{R}}$ introduced in the proof of Proposition 2.7. In fact, $h_\ell(t)$ acts on S^+ as multiplication by t^2 on $\mathbb{C}\ell$, by t^{-2} on $\mathbb{C}\bar{\ell}$ and it is trivial on $\langle \ell, \bar{\ell} \rangle^{\perp}$. The complex structure on $C(H)_{\mathbb{C}}^+ \cong V_{\mathbb{R}}^{\oplus 4}$, which defines the Kuga Satake variety $A_\ell \sim B_\ell^{\oplus 4}$, is also defined by h_ℓ ([18, Prop. 6.3]), now acting on $V_{\mathbb{R}}^4$. As $\rho_V(h_\ell) = h_{V,\ell}$, the complex structure is J_ℓ on $V_{\mathbb{R}}$. It follows that B_ℓ and $\mathcal{T}_\ell$ are isogeneous. $\square$

6.3. Remark

The proof of Proposition 6.2 uses the (algebraic) subgroup $Spin(H) = Spin_{h,s}$ of $Spin(S^+) = Spin(V)$. The decomposition $S_{\mathbb{Q}}^+ = H_{\mathbb{Q}} \oplus R_{\mathbb{Q}}$, with $R := \langle h, s \rangle$, implies that we actually have two commuting subgroups $Spin(H), Spin(R) \subset Spin(S^+)$.

Recall from Section 4.7 that $R_{\mathbb{C}} = \mathbb{C}\kappa \oplus \mathbb{C}\bar{\kappa}$ with $\kappa, \bar{\kappa} \in Q^+$. The decomposition of $V_{\mathbb{C}} = Z_{\kappa,\mathbb{C}} \oplus Z_{\bar{\kappa},\mathbb{C}}$ in the two isotropic eigenspaces for the K-action defines, as in Lemma 2.6, a one parameter subgroup h_R of $Spin(S_{\mathbb{R}}^+)$. As $h_R(t)\kappa = t^2\kappa$, $h_R(t)\bar{\kappa} = t^{-2}\bar{\kappa}$, this identifies the subgroup $Spin(R_{\mathbb{R}})$ with this one parameter subgroup, $h_R(U(1)) = Spin(R_{\mathbb{R}})$. In particular, the K-action on $V_{\mathbb{Q}}$ is generated by $Spin(R)$ and the scalar multiples of the identity.

The fact that $Spin(H), Spin(R) \subset Spin(S^+)$ commute implies that the subspaces $Z_{\kappa,\mathbb{C}}, Z_{\bar{\kappa},\mathbb{C}}$ are $Spin(H_{\mathbb{C}})$-invariant subspaces. Thus the spin representation of $Spin(H_{\mathbb{C}})$ on $V_{\mathbb{C}}$ is reducible. These two subspaces are the two half-spin representations of $Spin(H_{\mathbb{C}})$.

There is an isomorphism $Spin(H_{\mathbb{C}}) \cong SL(4, \mathbb{C})$ and the half-spin representations are identified with the standard representation $\mathbb{C}^4$ of $SL(4, \mathbb{C})$ and its dual $(\mathbb{C}^4)^*$. The

representation $H_{\mathbb{C}}$ is identified with $\wedge^2\mathbb{C}^4 \cong \wedge^2(\mathbb{C}^4)^*$, the isomorphism follows from the pairing, defined by the wedge product, $(\wedge^2\mathbb{C}^4) \times (\wedge^2\mathbb{C}^4) \to \wedge^4\mathbb{C}^4 \cong \mathbb{C}$.

6.4. The second cohomology group of $\mathcal{T}_\ell$

In [10] the Hodge structure on the second cohomology group $H^2(B, \mathbb{Q})$ of an abelian fourfold of Weil type with field K is studied. This group has dimension $\binom{8}{2} = 28$ and decomposes under the K-action into a $16 = 1 + 15$-dimensional subspace S'_B on which $x \in K$ acts as $x\bar{x}$, this subspace includes the polarization of Weil type. There is a complementary subspace S_B on which the eigenvalues of x are $x^2, \bar{x}^2$ of dimension 12. This subspace can be identified with the six dimensional K vector space $\wedge_K^2 H^1(B, K)$.

$$H^2(B, \mathbb{Q}) = S_B \oplus S'_B, \qquad S'_B := \{\xi \in H^2(B, \mathbb{Q}) : x^*\xi = x\bar{x}\xi, \quad \forall x \in K\}.$$

For a general fourfold of Weil type (so $SMT(B)_{\mathbb{R}} \cong SU(2,2)$) the Hodge structure S_B is a simple Hodge structure (so does not admit non-trivial Hodge substructures) if and only if the discriminant of B is non-trivial [10, Cor. 3.6].

In case the discriminant is trivial, one finds that $S_B \cong H_B^{\oplus 2}$, for a weight two, rank six, polarized, Hodge structure H_B which has Hodge numbers $(1, 4, 1)$. Moreover, the Kuga Satake variety of H_B is isogeneous to B^4, so one recovers the weight two Hodge structure H_B from its Kuga Satake variety.

The following proposition uses this result to show that the abelian fourfolds of Weil type $\mathcal{T}_\ell$ have trivial discriminant.

6.5. Proposition

For $\ell \in \Omega_{\{h,s\}^\perp}$, with h, s as in Theorem 4.6, the polarized abelian fourfold of Weil type $(\mathcal{T}_\ell, K, \omega_K)$ has trivial discriminant.

Proof. By [10, Cor. 3.6] it suffices to show that $(H_{\ell,\mathbb{Q}})^{\oplus 2}$ is isomorphic to the Hodge substructure $S_{\mathcal{T}_\ell} \subset H^2(\mathcal{T}_\ell, \mathbb{Q})$.

As in the proof of Proposition 2.7, the (weight one) Hodge structure on $V = H^1(\mathcal{T}_\ell, \mathbb{Z})$ defines a one parameter subgroup h_ℓ in $Spin(V)$ (actually in $Spin(V)_{h,s} \subset Spin(S^+) = Spin(V)$). A representation U of $Spin(V_{\mathbb{R}})$ on a real vector space U defines a Hodge decomposition $U_{\mathbb{C}} = \oplus U^{p,q}$, with $\overline{U^{p,q}} = U^{q,p}$, given by the eigenspaces $U^{p,q} = \{u \in U : h_\ell(z)u = z^a\bar{z}^b u\}$ (but the weight is not uniquely defined since $z\bar{z} = 1$).

The representation ρ^+ on $S_{\mathbb{R}}^+$ has the Hodge decomposition

$$(S^+)^{2,0} = H_\ell^{2,0} = \mathbb{C}\ell, \qquad (S^+)^{0,2} = \overline{(S^+)^{2,0}}, \qquad (S^+)^{1,1} = \left((S^+)^{2,0} \oplus (S^+)^{0,2}\right)^\perp$$

since these spaces are the eigenspaces for h_ℓ acting on $S_{\mathbb{C}}^+$ (see Lemma 2.6). The Hodge structure $S_{\mathbb{Q}}^+$ is a direct sum of Hodge structures

$$S_{\mathbb{Q}}^+ = H_{\ell,\mathbb{Q}} \oplus R_{\mathbb{Q}}, \qquad R := \langle h, s \rangle,$$

where $R_{\mathbb{Q}} \cong \mathbb{Q}(-1)^2$ is a trivial Hodge substructure with $R_{\mathbb{Q}}^{1,1} = R_{\mathbb{C}}$.

There is an isomorphism of $Spin(V) = Spin(S^+)$-representations $\wedge^2 S^+ = \wedge^2 V$ (both are the irreducible $so(8)$-representation with highest weight $(L_1 + L_2 + L_3 + L_4)/2 +$

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$(L_1 + L_2 - L_3 - L_4)/2 = L_1 + L_2)$. Hence we get a splitting of the Hodge structure on $\wedge^2 S_{\mathbb{Q}}^+$ (which is again defined by h_ℓ eigenspaces) in three Hodge substructures which have dimensions $\binom{6}{2} = 15$, $6 \cdot 2 = 12$ and 1 respectively:

$$\wedge^2 S_{\mathbb{Q}}^+ = (\wedge^2 H_{\ell,\mathbb{Q}}) \ \oplus \ (H_{\ell,\mathbb{Q}} \otimes R_{\mathbb{Q}}) \ \oplus \ (\wedge^2 R_{\mathbb{Q}}).$$

(The Hodge structure S^+ has weight two, so the Hodge structure on $\wedge^2 S^+$ should have weight four. However, $(\dim S^+)^{2,0} = 1$, so $\wedge^2 S_{\mathbb{Q}}^+$ has trivial $(4, 0)$ and $(0, 4)$ summands and thus it is the Tate twist of a weight two Hodge structure.)

Using the isomorphisms $\wedge^2 S_{\mathbb{Q}}^+ = \wedge^2 V = H^2(\mathcal{T}_\ell, \mathbb{Z})$ we see that

$$H_{\ell,\mathbb{Q}} \otimes R_{\mathbb{Q}} \ \cong \ (H_{\ell,\mathbb{Q}})^{\oplus 2} \ \hookrightarrow \ H^2(\mathcal{T}_\ell, \mathbb{Q})$$

is a non-simple Hodge substructure of $H^2(\mathcal{T}_\ell, \mathbb{Q})$.

It remains to check that $x \in K$ has eigenvalues $x^2, \bar{x}^2$ on this substructure. One can deduce this from the fact that representation $\wedge^2 V_{\mathbb{C}}$ of the complex Mumford Tate group $SL(4, \mathbb{C})$ of $\mathcal{T}_\ell$ is isomorphic to

$$\wedge^2(\mathbb{C}^4 \oplus (\mathbb{C}^4)^*) \cong (\wedge^2 \mathbb{C}^4)^{\oplus 2} \ \oplus \ \mathbb{C}^4 \otimes (\mathbb{C}^4)^*$$

and the last summand is the direct sum of a trivial one dimensional representation and an irreducible 15 dimensional representation. As the complexification of a Hodge substructure is a subrepresentation, there is a unique subrepresentation of dimension 12. Hence $S_{\mathcal{T}_\ell} = H_{\ell,\mathbb{Q}} \otimes R_{\mathbb{Q}}$ as desired.

Alternatively, by Remark 6.3, the $K^\times$-action is essentially given by the subgroup $Spin(R)$ of $Spin(S^+)$. This subgroup acts trivially on $\wedge^2 H_{\ell,\mathbb{Q}}$ and $\wedge^2 R_{\mathbb{Q}}$, so K acts through the norm on these summands. Therefore $S_{\mathcal{T}_\ell} = H_{\ell,\mathbb{Q}} \otimes R_{\mathbb{Q}} = (H_{\ell,\mathbb{Q}})^{\oplus 2}$. $\quad\square$

Declaration of competing interest

Data availability

No data was used for the research described in the article.

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